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THEORETICAL ANALYSIS OF TRANSONIC FLOW PAST UNSTAGGERED OSCILLA--ETC(U)
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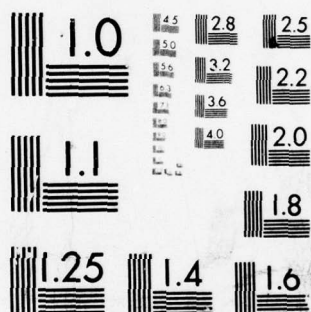
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THEORETICAL ANALYSIS OF
TRANSONIC FLOW PAST
UNSTAGGERED OSCILLATING CASCADES,

by

10 Peter Carlton/Olsen

11 Sept 1978

Thesis Advisor:

M.F. Platzter

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SECURITY CLASSIFICATION OF THIS PAGE (When Data Entered)

REPORT DOCUMENTATION PAGE		READ INSTRUCTIONS BEFORE COMPLETING FORM
1. REPORT NUMBER	2. GOVT ACCESSION NO.	3. RECIPIENT'S CATALOG NUMBER
4. TITLE (and Subtitle) Theoretical Analysis of Transonic Flow Past Unstaggered Oscillating Cascades		5. TYPE OF REPORT & PERIOD COVERED Engineer's Thesis September 1978
		6. PERFORMING ORG. REPORT NUMBER
7. AUTHOR(s) Peter Carlton Olsen		8. CONTRACT OR GRANT NUMBER(s)
9. PERFORMING ORGANIZATION NAME AND ADDRESS Naval Postgraduate School Monterey, California 93940		10. PROGRAM ELEMENT, PROJECT, TASK AREA & WORK UNIT NUMBERS
11. CONTROLLING OFFICE NAME AND ADDRESS Naval Postgraduate School Monterey, California 93940		12. REPORT DATE September 1978
		13. NUMBER OF PAGES 134
14. MONITORING AGENCY NAME & ADDRESS (if different from Controlling Office) Naval Postgraduate School Monterey, California 93940		15. SECURITY CLASS. (of this report) Unclassified
		15a. DECLASSIFICATION/DOWNGRADING SCHEDULE
16. DISTRIBUTION STATEMENT (of this Report) Approved for public release; distribution unlimited.		
17. DISTRIBUTION STATEMENT (of the abstract entered in Block 20, if different from Report)		
18. SUPPLEMENTARY NOTES		
19. KEY WORDS (Continue on reverse side if necessary and identify by block number) Oscillating Airfoil Transonic Flow Collocation Solution Potential Flow		
20. ABSTRACT (Continue on reverse side if necessary and identify by block number) → This paper presents an independent verification of the collocation method as a technique for calculating the lift on an oscillating airfoil in an unstaggered cascade immersed in transonic flow. This method was → next page		

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(20. ABSTRACT Continued)

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JUSTIFICATION	
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Theoretical Analysis of
Transonic Flow Past
Unstaggered Oscillating Cascades

by

Peter Carlton Olsen
Lieutenant, United States Coast Guard
B.S., United States Coast Guard Academy, 1970
M.S., University of West Florida, 1975
M.S.A.E., Naval Postgraduate School, 1977
M.S.O.R., Naval Postgraduate School, 1978

Submitted in partial fulfillment of the
requirements for the degree of

AERONAUTICAL ENGINEER

from the

NAVAL POSTGRADUATE SCHOOL

September 1978

Author

Peter C. Olsen

Approved by:

May F. Platter

Thesis Advisor

Robert E. Ball

Second Reader

May F. Platter

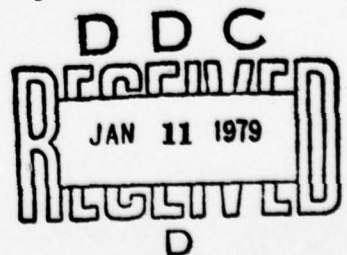
Chairman, Department of Aeronautics

William M. Toller

Dean of Science and Engineering

DISTRIBUTION STATEMENT A

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ABSTRACT

This paper presents an independent verification of the collocation method as a technique for calculating the lift on an oscillating airfoil in an unstaggered cascade immersed in transonic flow. This method was originally proposed by Gorelov. Results presented here differ somewhat from those presented by him. Two formulations are shown; one is purely numerical, the second employs an analytic expansion for small frequency.

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LIST OF SYMBOLS

a	= local speed of sound	II
a_o	= speed of sound in the uniform flow	III
c	= blade semichord	III
f_j	= elementary function used in collocation solution	V
F	= specific energy ("head")	II
G	= function describing the surface of the airfoil as a function of time	III
H	= function specifying location of blade surface in the vertical axis	III
i	= $\sqrt{-1}$	II, III, IV, V
k	= Strouhal number, nondimensional frequency	III, IV, V
m	= $\sqrt{(\gamma+1)w}$	IV, V
M	= Mach number = $\frac{ \vec{V} }{a}$	III
n	= number of collocation points - 1, order of highest spanning function	V
p	= pressure	II
	= nondimensional interblade distance	V
R	= universal gas constant	II
R.P.	= "real part of"	III, IV, V
T	= temperature	II
t	= time, nondimensional time	II, III, IV, V
U_o	= uniform velocity from infinity	II, III, VI
u	= x-component of velocity	II, III
u^o, u^1	= interference vertical velocities due to reference and adjacent blades respectively, solved so as to satisfy the tangential flow conditions	V, VI

u'	= small disturbance velocity	III
v	= y-component of velocity	II, III, IV, V
\vec{V}	= general velocity vector	II
v'	= small disturbance velocity	III
v^0, v^1	= vertical velocities due to the reference and adjacent blades respectively, determined from the tangential flow condition	V
w	= $\tilde{\phi}_x$, a constant used in Gorelov's approximation of the transonic flow potential	IV, V
x	= horizontal coordinate, may be non-dimensional	
x_*	= mp (transformed interblade distance in Gorelov's approximation)	V
x_ℓ	= blade leading edge	IV
x_0	= center of pitch of the unstaggered cascade	IV, V
y	= vertical coordinate, may be non-dimensional	
y, y_1	= vertical coordinates attached to the reference and adjacent blades respectively, may be non-dimensional	IV, V
z, z_1	= transformed vertical coordinates used in Gorelov's approximation, attached to the reference and adjacent blades respectively. $z = my$, $z_1 = my_1$	IV, V
$\frac{D}{Dt}$	= substantial derivative w.r.t. time $= \frac{\partial}{\partial t} + \frac{\partial x}{\partial t} \frac{\partial}{\partial x} + \frac{\partial y}{\partial t} \frac{\partial}{\partial y}$ $= \frac{\partial}{\partial t} + u \frac{\partial}{\partial x} + v \frac{\partial}{\partial y}$	II, III
$O(\omega^2)$	= "of the order of magnitude of ω^2 "	V
α	= angle of attack	II, III, IV, V
α_0	= maximum amplitude of pitch oscillations	IV

γ	= ratio of specific heats, c_p/c_v	
δ_{io}	= Dirac Delta function = 1 when $i = 0$ = 0 when $i \neq 0$	V
η	= $\cos^{-1}(-x)$	V
η_*	= $\cos^{-1}(1-x_*)$	V
$\bar{\eta}$	= $\cos^{-1}(-s)$	V
$\hat{\eta}$	= $\cos^{-1}(x_*-x)$	V
θ_j^0, θ_j^1	= interference potential coefficients for reference and adjacent blades respectively	V
λ	= k/m^2	IV,V
μ_j^0, μ_j^1	= Fourier coefficients describing the motion of the reference and adjacent blades respectively	V
ν	= angular frequency of oscillation	IV,V
ρ	= density	II
σ	= phase angle	V
τ	= cascade solidity, $\frac{2}{p}$	V
Φ	= general velocity potential	II,III,VI
Φ_0	= uniform flow velocity potential	
$\tilde{\Phi}$	= steady flow perturbation potential	III,IV
ϕ^0, ϕ^1	= perturbation potential in collocation solution due to reference and adjacent blades respectively	V,VI
ϕ^0, ϕ^1	= transformed potentials	V
ψ	= oscillatory flow potential	III,IV,V
Φ	= transformed oscillatory potential in Gorelov's coordinates corresponding to ψ	V
ψ^0, ψ^1	= interference potentials due to reference and adjacent blades respectively	V,VI

ψ^0, ψ^1 = transformed potentials in Gorelov's coordinates, corresponding to ψ^0 and ψ^1 V

$$\omega = \frac{k(1-m)^2}{m^4}$$

$\nabla = \vec{i} \frac{\partial}{\partial x} + \vec{j} \frac{\partial}{\partial y}$, Gradient operator, \vec{i} and \vec{j} are unit vectors in the x and y direction respectively

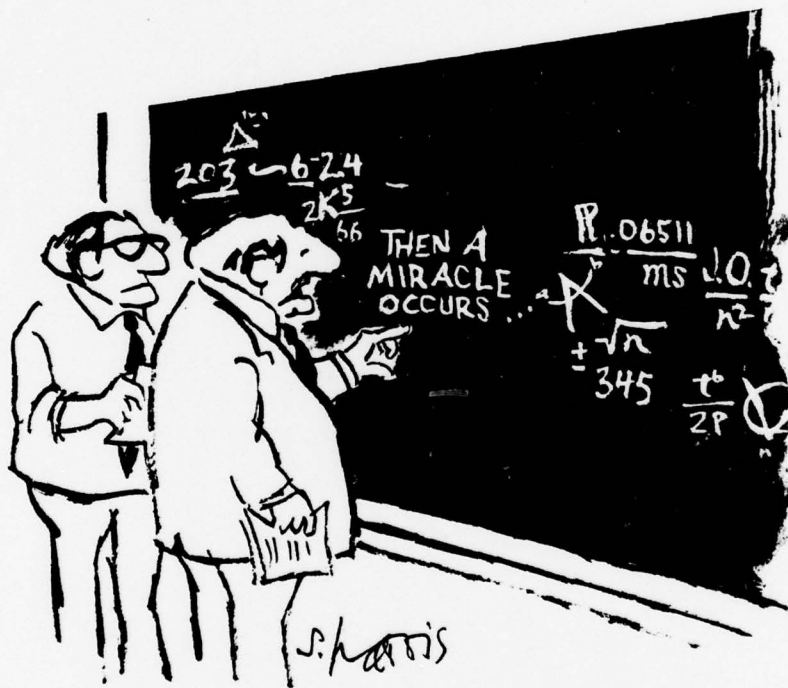
Computer Variables

DK	=	Reduced Frequency, k
DLAMDA	=	λ
DM2	=	m^2
DR	=	$mp = x_*$
ETA	=	η
ETASTR	=	η_*
IPT	=	Print Parameter
N	=	n
NF	=	not used
OFFSET	=	r
OMEGA	=	ω
QALPHA	=	$0 + i(\lambda - k)$
QCONST	=	$e^{i\sigma}$
QDCL	=	Cl_α
QDCM	=	Cm_α
QDK	=	$0 + ik$
QEXP	=	$0 + i\lambda$
QINTAP, QINTRP	=	Variables used to transmit boundary condition integrals
Q1ABCF, Q1RBCF	=	Interference coefficients for adjacent and reference blades
Q1COF	=	Right hand side vector in collocation solution
Q1INT	=	Known matrix of integrals in collocation solution
Q1RBP, Q1ABP	=	Not used
Q2CP	=	Not used
Q2EXP	=	$e^{-i\lambda x}$

Q2PT	=	Not used
RHO	=	vertical displacement
RHP	=	local input variable for rho
SIGMA	=	σ
TAU	=	τ
XASTN	=	Current x station in adjacent blade coordinates
XSTN	=	Current x station in reference blade coordinates

ACKNOWLEDGEMENT

I would like to thank Professor Max F. Platzer without whose infinite patience this thesis would not have been written.



"I think you should be more explicit here in step two."

Cartoon Reprinted from American Scientist, Vol. 65, No. 6, Nov-Dec 1977 with permission from Sidney Harris.

I. INTRODUCTION

The analysis of unsteady transonic flows in aircraft turbopropulsion is an area of intense current interest. Rising fuel prices and increasing thrust requirements both point toward the need of turbomachinery capable of performing well with transonic or supersonic internal flow. But, increased flow has increased both the costs and uncertainties of engine designs. Flutter problems have already become a major consideration in engine development. Problems unforeseen in earlier days of turbine engine production have caused long development delays, or forced acceptance of engines producing less than their initial design thrust. These uncertainties cannot be avoided when an attempt is made to extend the state of the art, but they can be reduced by extending the range of analytical modeling.

Such extension must now be done piecemeal. The three-dimensional flows in turbomachinery, including the simultaneous effects of boundary layers, rotation, finite blade thickness, spanwise Mach distributions, and shocks, are well beyond present capability. Perhaps one day complete analysis will be practical, but it is not today. The best that can be done now is to approach the problem from one aspect at a time. Flow through a two dimensional cascade has been a useful tool in this partial analysis.

This thesis was originally to have been an extension of the work of Elder [1] and Schlein [2] to the case of a staggered cascade. Their work, based on Teipel's [3] linearization of the unsteady transonic small perturbation equation, analyzed transonic flow through oscillating unstaggered cascades by use of the collocation method. While the problem was easy to state, it was difficult to solve. Both Elder and Schlein had encountered difficulty in employing the collocation method. Therefore, it was decided that verification of the basic collocation solution presented by Gorelov [4] using a different linearization would be a worthwhile goal in itself.

The following investigation presents a verification of the development in [4], along with numerical results and suggestions for further work.

II. UNSTEADY TRANSONIC FLOW THEORY

Considering inviscid flow only, the following four equations govern the aerodynamic flow problem at hand:

The equation of state

$$p = \rho RT \quad (\text{II-1})$$

and the equations for the conservation of

$$1. \text{ Mass: } \operatorname{div}(\rho \vec{v}) + \frac{\partial \rho}{\partial t} = 0 \quad (\text{II-2})$$

$$2. \text{ Momentum: } \frac{D\vec{v}}{Dt} + \frac{1}{\rho} \nabla p = 0 \quad (\text{II-3})$$

$$3. \text{ Energy: } \frac{DS}{Dt} = 0 \quad (\text{II-4})$$

where

\vec{v} = velocity

p = pressure

S = entropy

R = universal gas constant

T = temperature

t = time

ρ = density

$$\frac{D}{Dt} = \frac{\partial}{\partial t} + \frac{\partial}{\partial x} \frac{\partial x}{\partial t} + \frac{\partial}{\partial y} \frac{\partial y}{\partial t}$$

The analysis starts with a uniform flow from infinity. This flow has velocity U_0 parallel to the x-axis. This

formulation can be simplified by working with the total velocity potential, ϕ , where

$$u = \frac{\partial \phi}{\partial x} = \phi_x = x \text{ component of velocity} = \frac{\partial x}{\partial t} \quad (\text{II-5})$$

$$v = \frac{\partial \phi}{\partial y} = \phi_y = y \text{ component of velocity} = \frac{\partial y}{\partial t} \quad (\text{II-6})$$

Thus, the initial uniform flow is represented by the uniform flow potential

$$\phi_o = U_o x \quad (\text{II-7})$$

This notation may be applied to the conservation equations for mass and momentum. The equation for conservation of mass

$$\text{div}(\rho \vec{v}) + \frac{\partial \rho}{\partial t} = 0 \quad (\text{II-8})$$

becomes for two-dimensional unsteady flow

$$\begin{aligned} \frac{\partial(\rho u)}{\partial x} + \frac{\partial(\rho v)}{\partial y} + \frac{\partial \rho}{\partial t} &= 0 \\ \left[\frac{\partial \rho}{\partial t} + u \frac{\partial \rho}{\partial x} + v \frac{\partial \rho}{\partial y} \right] + \rho \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) &= 0 \end{aligned} \quad (\text{II-9})$$

but

$$\frac{\partial \rho}{\partial t} + u \frac{\partial \rho}{\partial x} + v \frac{\partial \rho}{\partial y} = \frac{D\rho}{Dt}$$

and

$$u = \phi_x \quad \text{and} \quad v = \phi_y$$

Thus

$$\frac{D\rho}{Dt} + \rho(\phi_{xx} + \phi_{yy}) = 0$$

and

$$\phi_{xx} + \phi_{yy} = -\frac{1}{\rho} \frac{D\rho}{Dt} \quad (\text{II-10})$$

The speed of sound is given by

$$a^2 = \frac{dp}{d\rho}$$

Thus

$$\frac{D\rho}{Dt} = \frac{d\rho}{dp} \cdot \frac{Dp}{Dt} = \frac{1}{a^2} \frac{Dp}{Dt}$$

Applying this to equation (II-10) yields

$$\phi_{xx} + \phi_{yy} = -\frac{1}{\rho a^2} \frac{Dp}{Dt} \quad (\text{II-11a})$$

$$\phi_{xx} + \phi_{yy} = -\frac{1}{\rho a^2} (u p_x + v p_y + p_t) \quad (\text{II-11b})$$

$$= -\frac{1}{\rho a^2} [(\nabla\phi) \cdot (\nabla p) + p_t] \quad (\text{II-11c})$$

where ∇ is the gradient operator, $P_x = \frac{\partial P}{\partial x}$, $P_y = \frac{\partial P}{\partial y}$, $P_t = \frac{\partial P}{\partial t}$.

Laying this aside for the moment, consider the momentum equation (II-3)

$$\frac{D\vec{v}}{Dt} + \frac{1}{\rho} \nabla p = 0$$

$$\frac{D\vec{v}}{Dt} = \frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \nabla) \vec{v}$$

Thus

$$\rho \left[\frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \nabla) \vec{v} \right] = \nabla p \quad (\text{II-12})$$

$$\vec{v} = \nabla \phi$$

so

$$\frac{\partial \vec{v}}{\partial t} = \frac{\partial}{\partial t} (\nabla \phi) = \nabla \frac{\partial \phi}{\partial t} \quad (\text{II-13})$$

and

$$(\vec{v} \cdot \nabla) \vec{v} = \nabla \frac{v^2}{2} - \vec{v} \times (\nabla \times \vec{v})$$

where

$$v^2 = u^2 + v^2$$

$$= \vec{v} \cdot \vec{v}$$

For irrotational flow

$$\nabla \times \vec{v} = 0$$

thus

$$(\vec{v} \cdot \nabla) \vec{v} = \frac{\nabla v^2}{2} \quad (\text{II-14})$$

Thus

$$\frac{\nabla p}{\rho} + \nabla \left[\phi_t + \frac{v^2}{2} \right] = 0 \quad (\text{II-15})$$

which after integration along a streamline becomes

$$\int \frac{dp}{\rho} + \phi_t + \frac{v^2}{2} = F(t) \quad (\text{II-16})$$

For uniform flow from infinity $F(t) = \frac{1}{2} U_o^2$ and thus the final result is

$$\int \frac{dp}{\rho} + \phi_t + \frac{v^2}{2} = \frac{1}{2} U_o^2 \quad (\text{II-17})$$

Differentiation with respect to t gives

$$p_t = -\rho(\phi_{tt} + \frac{1}{2} \frac{\partial V^2}{\partial t}) \quad (\text{II-18})$$

From (II-3)

$$-\nabla p = \rho \frac{D\vec{v}}{Dt} \quad (\text{II-19})$$

Substitute (II-18) and (II-19) into (II-11c) to obtain

$$\phi_{xx} + \phi_{yy} = \frac{1}{a^2} [(\nabla\phi) \cdot \frac{D\vec{v}}{Dt} + \phi_{tt} + \frac{1}{2} \frac{\partial V^2}{\partial t}] \quad (\text{II-20})$$

This may be further simplified

$$\begin{aligned} \frac{D\vec{v}}{Dt} &= \frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \nabla) \vec{v} \\ &= \frac{\partial \vec{v}}{\partial t} + \frac{1}{2} \nabla V^2 \end{aligned}$$

Hence

$$\phi_{xx} + \phi_{yy} = \frac{1}{a^2} [(\nabla\phi) \cdot (\frac{\partial \vec{v}}{\partial t} + \frac{1}{2} \nabla V^2) + \phi_{tt} + \frac{1}{2} \frac{\partial V^2}{\partial t}] \quad (\text{II-21})$$

Expanding terms

$$\begin{aligned} (\nabla\phi) \cdot (\frac{\partial \vec{v}}{\partial t}) &= \nabla\phi \cdot (\frac{\partial}{\partial t} \nabla\phi) = \phi_x \phi_{xt} + \phi_y \phi_{yt} \\ \nabla\phi \cdot \frac{\nabla V^2}{2} &= \frac{\phi_x^2 \phi_{xx}}{2} + \frac{\phi_y^2 \phi_{yy}}{2} + \frac{\phi_x \phi_y \phi_{xy}}{2} + \frac{\phi_y \phi_x \phi_{yx}}{2} \\ &= \frac{\phi_x^2 \phi_{xx}}{2} + \frac{\phi_y^2 \phi_{yy}}{2} + \frac{2\phi_x \phi_y \phi_{xy}}{2} . \end{aligned}$$

$$\begin{aligned}
\frac{1}{2} \frac{\partial v^2}{\partial t} &= \frac{1}{2} \frac{\partial}{\partial t} [(\nabla \phi) \cdot (\nabla \phi)] \\
&= \frac{1}{2} [\phi_x \phi_{xt} + \phi_{xt} \phi_x + \phi_y \phi_{yt} + \phi_{yt} \phi_y] \\
&= \phi_x \phi_{xt} + \phi_y \phi_{yt}
\end{aligned}$$

The final result obtained by combining terms is

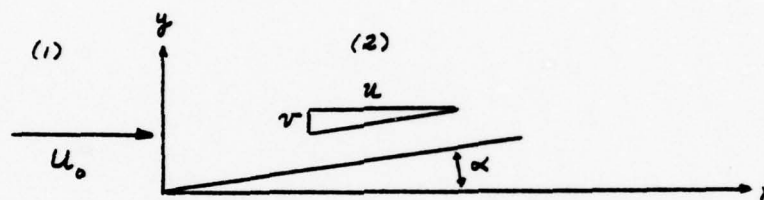
$$\begin{aligned}
\left(1 - \frac{\phi_x^2}{a^2}\right)_{xx} + \left(1 - \frac{\phi_y^2}{a^2}\right)_{yy} - \frac{2\phi_x \phi_y \phi_{xy}}{a^2} \\
- \frac{2\phi_x}{a^2} \phi_{xt} - \frac{2\phi_y}{a^2} \phi_{yt} - \frac{1}{a^2} \phi_{tt} = 0 \quad (\text{II-22})
\end{aligned}$$

This result is valid for irrotational, inviscid, two-dimensional, unsteady, compressible flows where gravity has been neglected.

III. SMALL PERTURBATION THEORY OF TRANSONIC FLOW

A. GENERAL CASE

A thin body at a small angle of attack will cause only a slight disturbance in the fluid. A flat plate is an example. Consider flow past a flat plate at angle of attack, α .



The flow at (2) must be parallel to the plate. To achieve this, small disturbance velocities u' and v' must be added to the free stream velocity yielding

$$u = U_0 + u'$$

$$v = v'$$

The potential of the disturbed flow may be considered as the sum of the uniform flow potential, $\phi_0 = U_0 x$, and the disturbance potential, ϕ

$$\phi = \phi_0 + \phi \quad (\text{III-1})$$

Thus

$$\phi_x = U_0 + u' \quad (\text{III-2a})$$

$$\phi_y = v \quad (\text{III-2b})$$

If ϕ is a function of time, then

$$\phi_t = \phi_t$$

This result may be substituted into (II-22) leading to

$$\begin{aligned} & \left[1 - \frac{(U_0 + u')^2}{a^2}\right] \phi_{xx} + \left[1 - \frac{v^2}{a^2}\right] \phi_{yy} - 2 \frac{(U_0 + u')v}{a^2} \phi_{xy} \\ & - 2 \frac{(U_0 + u')}{a^2} \phi_{xt} - \frac{2v}{a^2} \phi_{yt} - \frac{1}{a^2} \phi_{tt} = 0 \quad (\text{III-3}) \end{aligned}$$

This expands to yield

$$\begin{aligned} & \left[1 - \frac{U_0^2 + 2U_0u' + u'^2}{a^2}\right] \phi_{xx} + \left[1 - \frac{v^2}{a^2}\right] \phi_{yy} - 2 \frac{(U_0v + u'v)}{a^2} \phi_{xy} \\ & - 2 \frac{(U_0 + u')}{a^2} \phi_{xt} - \frac{2v}{a^2} \phi_{yt} - \frac{1}{a^2} \phi_{tt} = 0 \quad (\text{III-4}) \end{aligned}$$

Equation (III-4) may be further simplified as shown by Landahl [3]. All non-linear terms except the $\phi_x \phi_{xx}$ product

term can be neglected yielding the following transonic small disturbance equation

$$\begin{aligned} & [(M^2-1) + (\gamma+1)M^2 \frac{\phi_x}{U_0}] \phi_{xx} \\ & - \phi_{yy} + \frac{2M_0}{a_0} \phi_{xt} + \frac{1}{a_0} \phi_{tt} = 0 \end{aligned} \quad (\text{III-5})$$

where a_0 is the velocity of sound in the free-stream, and γ is the ratio of specific heats, and M_0 is the Mach number.

B. BOUNDARY CONDITION

The tangential flow condition requires that the flow be tangent to the airfoil surface at each instant of time. This means that no fluid may flow through the surface of the airfoil and is expressed by the condition

$$\frac{DG}{Dt} = 0 \quad \text{on } G(x,y,t) \quad (\text{III-6})$$

where

$G(x,y,t)$ describes the surface of the body as a function of time.

For a thin airfoil restricted to small oscillations, this may be written as

$$G = y - H(x,t) \quad (\text{III-7})$$

where

$H(x,t)$ is the function describing the position of the airfoil.

$H(x,t)$ can be written for harmonic pitch oscillations as

$$H(x,t) = \text{R.P.}[\alpha_0(x-x_0) e^{i\nu t}] \quad (\text{III-8})$$

where the time-varying angle of attack $\alpha(t)$ is given by

$$\alpha(t) = \text{R.P.}[\alpha_0 e^{i\nu t}]$$

and α_0 = maximum amplitude of pitch oscillation

x_0 is the pitch axis

ν is the angular frequency of oscillation

$$i = \sqrt{-1}$$

R.P. = "real part of"

Inserting (III-8) into the flow tangency condition (III-6) gives, after linearization,

$$\phi_y(x,0) = v(x,0) = \alpha_0[U_0 + i\nu(x - x_0)] e^{i\nu t} \quad (\text{III-9})$$

on $y = 0$

This is a condition for the normal velocity to be prescribed at the airfoil's mean position $y = 0$.

C. NONDIMENSIONALIZATION

The terms in equations (III-5) and (III-9) are dimensional. For the following calculations it is convenient to use non-dimensional quantities. Define non-dimensional time and length to be

$$\bar{x} = \frac{x}{c}$$

$$\bar{y} = \frac{y}{c} \quad (\text{III-10})$$

$$\bar{t} = \frac{tU_o}{c}$$

where

U_o = uniform velocity from infinity

c = reference length (blade semichord).

The velocity potential in equation (III-5) may be non-dimensionalized as follows. Let

$$\bar{\phi} = \frac{\phi}{U_o c}$$

Hence:

$$\phi = U_o c \bar{\phi}$$

$$\phi_x = U_o c \bar{\phi}_x \left(\frac{1}{c}\right)$$

$$= U_o \bar{\phi}_x \quad (\text{III-11})$$

and similarly for the other derivatives in (III-5), yielding

$$[(M^2-1) + (\gamma+1)M^2\bar{\phi}_{\bar{x}}]\bar{\phi}_{\bar{x}\bar{x}} - \bar{\phi}_{\bar{y}\bar{y}} + 2M^2\bar{\phi}_{\bar{x}\bar{t}} + M^2\bar{\phi}_{\bar{t}\bar{t}} = 0$$

(III-12)

This equation is non-dimensional.

The boundary condition given in equation (III-9) may be non-dimensionalized in a similar fashion

$$\phi_y(x,0) = \alpha_o [U_o + i\nu(x - x_o)] e^{i\nu t} \quad (\text{III-9})$$

Thus

$$U_o c \bar{\phi}_y = \alpha_o [U_o + ik \frac{U_o}{c} (c\bar{x} - c\bar{x}_o)] e^{i\nu \bar{t} \frac{c}{U_o}} \quad (\text{III-13})$$

where

$$k = \frac{\nu c}{U_o} = \text{Strouhal number or reduced frequency}$$

$$U_o c \bar{\phi}_y \cdot \frac{1}{c} = \alpha_o U_o [1 + ik(\bar{x} - \bar{x}_o)] e^{ik\bar{t}} \quad (\text{III-14})$$

Thus

$$\bar{\phi}_y = \bar{v}(\bar{x},0) = \alpha_o [1 + ik(\bar{x} - \bar{x}_o)] e^{ik\bar{t}} \quad (\text{III-15})$$

Because the final operations are linear in α_o , set $\alpha_o = 1$, yielding

$$\bar{\phi}_y = \bar{v}(\bar{x},0) = [1 + ik(\bar{x} - \bar{x}_o)] e^{ik\bar{t}} \quad (\text{III-16})$$

The overbars denoting nondimensional quantities will be dropped from the remainder of the paper. All further quantities shall be assumed appropriately non-dimensional. This yields the following final equations

$$[(M^2-1)+(\gamma+1)M^2\phi_x]\phi_{xx} - \phi_{yy} + 2M^2\phi_{xt} + M^2\phi_{tt} = 0 \quad (\text{III-17})$$

and

$$\phi_y(x,0) = v(x,0) = [1 + ik(x - x_0)] e^{ikt} \quad (\text{III-18})$$

where

all quantities are nondimensional and

$$\alpha_0 = 1$$

D. HARMONIC OSCILLATIONS

In the case of harmonic oscillations, equation (III-17) may be simplified still further.

Let

$$\phi = \tilde{\phi} + \text{R.P.}[\psi e^{ikt}]$$

where

$\tilde{\phi}$ = non-dimensional steady flow potential

ψ = non-dimensional oscillatory flow potential

R.P. = "real part of"

Equation (III-17) then becomes

$$\begin{aligned} (1-M^2)\psi_{xx} + \psi_{yy} - M^2(\gamma+1)[\psi_x\psi_{xx} + \tilde{\phi}_x\psi_{xx} + \tilde{\phi}_{xx}\psi_x] \\ + M^2k^2\psi - 2iMk^2\psi_x = 0 \end{aligned} \quad (\text{III-19})$$

For M close to 1, this is a nonlinear mixed elliptic-hyperbolic partial differential equation with variable coefficients, the exact type depending on $\tilde{\phi}_x$ and $\tilde{\phi}_{xx}$. However, because flutter analysis is primarily concerned with the stability of small perturbations about a steady flow, the oscillatory component may be assumed small compared to the steady flow potential and therefore the product term $\psi_x\psi_{xx}$ may be neglected, yielding,

$$\begin{aligned} (1-M^2)\psi_{xx} + \psi_{yy} - M^2(\gamma+1)[\tilde{\phi}_x\psi_{xx} + \tilde{\phi}_{xx}\psi_x] \\ = 2iM^2k\psi_x + M^2k^2\psi = 0 \end{aligned} \quad (\text{III-20})$$

IV. LINEARIZATION OF THE GOVERNING EQUATION

The basic flutter equation, (III-20), is still a non-linear, mixed elliptic-hyperbolic partial differential equation with variable coefficients and difficult to solve. It may yet be further simplified.

A. BASIC SOLUTION

For $M = 1$, equation (III-20) becomes

$$\psi_{YY} - (\gamma+1) [\tilde{\phi}_x \psi_{xx} + \tilde{\phi}_{xx} \psi_x] - 2ik\psi_x + k^2\psi = 0 \quad (\text{IV-1})$$

Now assume

$$\tilde{\phi}_x \approx w = \text{constant} \quad (\text{IV-2})$$

$$\tilde{\phi}_{xx} \approx 0$$

throughout the interblade channel. Setting

$$\tilde{\phi}_x (\gamma+1) = w(\gamma+1) = m^2 \quad (\text{IV-3})$$

yields

$$m^2 \psi_{xx} - \psi_{YY} + 2ik\psi_x - k^2\psi = 0 \quad (\text{IV-4})$$

The solution to this equation is found in Garrick and Rubinow [5]

$$\psi(x,y) = -\frac{1}{m} \int_{x_l}^{x-my} v(s) J_0[\omega \sqrt{(x-s)^2 - (my)^2}] e^{i\lambda(s-x)} ds \quad \text{for } y > 0 \quad (\text{IV-5a})$$

and

$$\psi(x,y) = \frac{1}{m} \int_{x_l}^{x+my} v(s) J_0[\omega \sqrt{(x-s)^2 - (my)^2}] e^{i\lambda(s-x)} dx \quad \text{for } y < 0 \quad (\text{IV-5b})$$

where

$$v(x) = \lim_{y \rightarrow 0} \frac{\partial}{\partial y} \psi(x,y) .$$

$v(x)$ may be obtained directly from the tangential flow boundary condition, and

$$\omega = \frac{k^2(1-m^2)}{m^4}$$

$$\lambda = \frac{k}{m^2} \sqrt{1+m^2} \approx \frac{k}{m^2} \quad \text{(where this paper employs the approximation used by Gorelov [4])}$$

$$x_l = \text{blade leading edge,}$$

Gorelov [4], has proposed a further simplification.

Set

$$z = my \quad (IV-6a)$$

$$\Psi(x, z) = \psi(x, y) e^{i\lambda x} \quad (IV-6b)$$

Equation (IV-4) then becomes

$$\Psi_{xx} - \Psi_{zz} + \omega^2 \Psi = 0$$

with solution

$$\Psi(x, z) = -\frac{1}{m} \int_{x_\ell}^{x-z} v(s) J_0[\omega \sqrt{(x-s)^2 - z^2}] e^{i\lambda s} ds \quad (IV-7a)$$

$$z > 0$$

$$\Psi(x, z) = \frac{1}{m} \int_{x_\ell}^{x-z} v(s) J_0[\omega \sqrt{(x-s)^2 - z^2}] e^{i\lambda s} ds \quad (IV-7b)$$

$$z < 0$$

where

$$v(x) = m e^{-i\lambda x} \lim_{z \rightarrow 0} \Psi_z(x, z)$$

$v(x)$ is obtained from the tangential flow boundary condition.

For a thin body immersed in the flow, the solutions for $y > 0$, $z > 0$, and $y < 0$, $z < 0$ apply above the body along left-running Mach lines, or below along right running Mach lines respectively.

B. BOUNDARY CONDITIONS

1. Flow Tangency Condition

The boundary condition comes from the tangential flow condition, (III-18)

$$v(x) = [1 + ik(x - x_0)] \quad (IV-8)$$

2. Upstream Condition

The final linearized equation is a hyperbolic differential equation with boundary condition

$$\psi(x, y) = 0 \quad (IV-9)$$

when

$$x - x_\ell < |my|$$

for the solution shown in equations (IV-5) or

$$\Psi(x, z) = 0$$

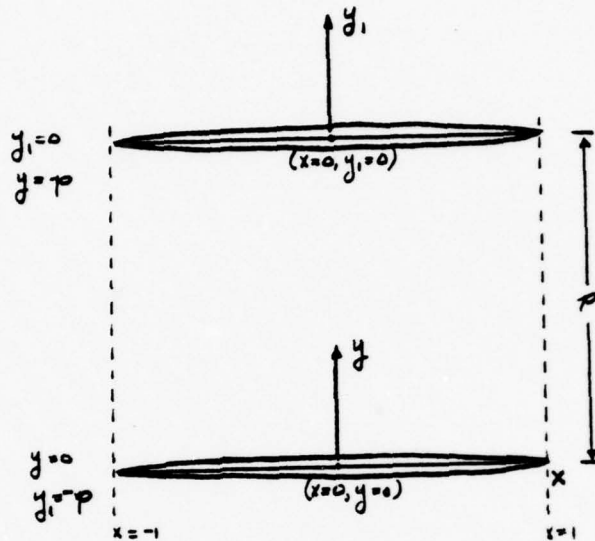
when

$$x - x_{\ell} < |z|$$

for the solution shown in equations (IV-7).

V. PROBLEM FORMULATION

A. CO-ORDINATE SYSTEM



Assume the geometry shown above. Both blades are thin airfoils of semichord c . All measurements are non-dimensional, normalized to c . The (x, y) co-ordinate system has its origin at the center of the reference (lower) blade. The (x, y_1) system is centered at the middle of the adjacent (upper) blade. The origin of the (x, y_1) system is located at $(0, p)$ in the reference system. Generalizing this convention, the same symbols shall be used for the same quantities on both blades. Where discrimination is required, the quantity associated with the adjacent blade will be marked with

superscript ¹, the quantity associated with the reference blade will be either unsuperscripted or marked with a superscript ⁰.

Each blade is assumed to perform a small amplitude harmonic oscillation about its mid-chord point. Both blades are assumed to have identical reduced frequencies, k , and the motion of the adjacent blade lags that of the reference blade by a phase angle σ .

The blades are immersed in a uniform flow from the left at $M = 1$. The objective is to determine the oscillatory pressure distributions and aerodynamic forces generated by the blades' oscillations. Cascade solidity, $\tau = 2/p$.

B. BOUNDARY CONDITIONS

1. Upstream Condition

$$\begin{aligned} \psi &= 0 && \text{whenever} \\ x + 1 &< |my| && (V-1) \\ &\text{and} \\ x + 1 &< |my_1|, && \text{simultaneously} \end{aligned}$$

2. Flow Tangency Condition

Along the reference blade

$$\lim_{y \rightarrow 0} \psi_y(x, y) = (1 + ikx) \quad (V-2a)$$

Along the adjacent blade

$$\lim_{y_1 \rightarrow 0} \psi_{y_1}(x, y_1) = (1 + ikx)e^{i\sigma} \quad (V-2b)$$

where

σ is the phase angle between the blades oscillations

C. BASIC SOLUTION TECHNIQUE

Assume that the unsteady potential, ψ , may be written as the sum of four components

$$\psi(x, y) = \phi^0(x, y) + \psi^0(x, y) + \phi^1(x, y_1) + \psi^1(x, y_1) \quad (V-3)$$

where:

- ϕ^0 = potential due to the reference blade alone, known from equation (IV-7)
- ϕ^1 = potential due to the adjacent blade alone, known from equation (IV-7)
- ψ^0 = interference potential required to satisfy tangential flow condition along reference blade, unknown
- ψ^1 = interference potential required to satisfy tangential flow condition along adjacent blade, unknown.

This total potential must satisfy the tangential flow condition at the plane of both the reference and adjacent blades. Thus

$$\begin{aligned} \phi_Y^0(x, y=0) + \phi_{Y_1}^1(x, y_1=-p) + \psi_Y^0(x, y=0) + \psi_{Y_1}^1(x, y_1=-p) \\ = (1 + ikx) \end{aligned} \quad (V-4a)$$

at the reference blade, and

$$\begin{aligned} \phi^0(x, y=p) + \phi_{y_1}^1(x, y_1=0) + \psi_Y^0(x, y=p) + \psi_{y_1}^1(x, y_1=0) \\ = (1 + ikx) e^{i\sigma} \end{aligned} \quad (V-4b)$$

at the adjacent blade.

But from the unsteady potential solution for a single oscillating blade one has

$$\phi_Y^0(x, y=0) = 1 + ikx \quad (V-5a)$$

and

$$\phi_{y_1}^1(x, y_1=0) = (1 + ikx) e^{i\sigma} \quad (V-5b)$$

Thus

$$\phi_{y_1}^1(x, y_1=-p) + \psi_{y_1}^1(x, y_1=-p) + \psi_Y^0(x, y=0) = 0 \quad (V-6a)$$

along the reference blade, and

$$\phi_Y^0(x, y=p) + \psi_Y^0(x, y=p) + \psi_{y_1}^1(x, y_1=0) = 0 \quad (V-6b)$$

along the adjacent blade.

From equation (IV-7)

$$\phi^0(x, y) = -\frac{1}{m} \int_{-1}^{x-my} v^0(s) J_0[\omega \sqrt{(x-s)^2 - (my)^2}] e^{i\lambda(s-x)} ds$$

$$y > 0 \quad (V-7a)$$

$$= \frac{1}{m} \int_{-1}^{x+my} v^0(s) J_0[\omega \sqrt{(x-s)^2 - (my)^2}] e^{i\lambda(s-x)} ds$$

$$y < 0 \quad (V-7b)$$

where

$$\begin{aligned} v^0(s) &= 1 + iks \\ \lambda &= k/m^2 \\ m &= (\gamma+1)w \\ \omega &= \frac{k^2(1-m^2)}{m^4} \\ w &= \text{mean value of } \tilde{\phi}_x \text{ in the channel} \end{aligned}$$

$$\phi^1(x, y_1) = -\frac{1}{m} \int_{-1}^{x-my_1} v^1(s) J_0[\omega \sqrt{(x-s)^2 - (my_1)^2}] e^{i\lambda(s-x)} ds$$

$$y_1 > 0 \quad (IV-8a)$$

$$= \frac{1}{m} \int_{-1}^{x+my_1} v^1(s) J_0[\omega \sqrt{(x-s)^2 - (my_1)^2}] e^{i\lambda(s-x)} ds$$

$$y_1 < 0 \quad (IV-8b)$$

where

$$v^1(s) = (1 + iks)e^{i\sigma}$$

Henceforth attention will be restricted to the flow within the channel, $0 \leq y \leq p$, $-p \leq y_1 \leq 0$ leaving (IV-7a) and (IV-8b) as the governing equations of interest.

The two interference potentials are assumed to have forms identical to (IV-7a) and (IV-8b).

Set

$$\psi^0(x, y) = -\frac{1}{m} \int_{-1}^{x-my} u^0(s) J_0[\omega \sqrt{(x-s)^2 - (my)^2}] e^{i\lambda(s-x)} ds$$

$y > 0$ (V-9a)

$$\psi^1(x, y_1) = \frac{1}{m} \int_{-1}^{x+m} u^1(s) J_0[\omega \sqrt{(x-s)^2 - (my_1)^2}] e^{i\lambda(s-x)} ds$$

$y_1 < 0$ (V-9b)

where

$u^0(s)$ and $u^1(s)$ are unknown functions to be determined so as to satisfy equations (V-6)

Substitution of (V-7a), (V-8b) and (V-9) into (V-6) yields

$$u^0(x) + \psi_{y_1}^1(x, y_1 = -p) + \phi_{y_1}^0(x, y_1 = -p) = 0 \quad (V-10a)$$

$$u^1(x) + \psi_Y^0(x, y=p) + \phi_Y^0(x, y=p) = 0 \quad (V-10b)$$

Recalling Gorelov's transformation discussed in [4] and shown in equations (IV-6) above, set

$$\phi^0(x, z) = \phi^0(x, y) e^{i\lambda x}$$

$$\phi^1(x, z_1) = \phi^1(x, y_1) e^{i\lambda x}$$

$$\psi^0(x, z) = \psi^0(x, y) e^{i\lambda x}$$

$$\psi^1(x, z_1) = \psi^1(x, y_1) e^{i\lambda x}$$

where $z = my$

$$z_1 = my_1$$

Then

$$\frac{e^{i\lambda x}}{m} u^0(x) + \phi_{z_1}^1(x, z_1 = -x_*) + \psi_{z_1}^1(x, z_1 = -x_*) = 0 \quad (V-11a)$$

and

$$\frac{e^{i\lambda x}}{m} u^1(x) + \phi_z^0(x, z = x_*) + \psi_z^0(x, z = x_*) = 0 \quad (V-11b)$$

where

$$x_* = mp.$$

To employ the collocation method, assume that $u^1(x)$ and $u^0(x)$ can be approximated as the sum of a set of elementary functions f_j so that

$$u^0(x) \approx \sum_{j=0}^n \theta_j^0 f_j(x) \quad (V-12a)$$

$$u^1(x) \approx \sum_{j=0}^n \theta_j^1 f_j(x) \quad (V-12b)$$

where $f_j(x) = 0$ when $x \leq x_* - 1$

Note that here both u^0 and u^1 are expressed in terms of the same elementary functions, f_j .

Because of the slightly supersonic nature of the problem observe that $u^0(x) = 0$ and $u^1(x) = 0$ when $x \leq x_* - 1$.

Equations (V-12) may now be rewritten as

$$e^{i\lambda x} \sum_{j=0}^n \theta_j^0 f_j(x) + \frac{\partial}{\partial z_1} \int_{x_*-1}^{x+z_1} \sum_{j=0}^n \theta_j^1 f_j(s) J_0[\omega \sqrt{(x-s)^2 - z_1^2}] e^{i\lambda s} ds$$

$$= -e^{i\sigma} \frac{\partial}{\partial z_1} \int_{-1}^{x+z_1} (1+iks) J_0[\omega \sqrt{(x-s)^2 - z_1^2}] e^{i\lambda s} ds$$

$$\text{at } z_1 = -x_* \quad (V-13a)$$

$$\begin{aligned}
e^{i\lambda x} \sum_j^1 f_j(x) - \frac{\partial}{\partial z} \int_{x_*-1}^{x-z} \sum_j^0 f_j(s) J_0[\omega \sqrt{(x-s)^2 - z^2}] e^{i\lambda s} ds \\
= \frac{\partial}{\partial z} \int_{-1}^{x+z} (1+iks) J_0[\omega \sqrt{(x-s)^2 - z^2}] e^{i\lambda s} ds
\end{aligned}$$

$$\text{at } z = x_* \quad (V-13b)$$

where

$$f_j(x) = 0 \text{ for } x \leq x_* - 1$$

This simplifies to

$$\begin{aligned}
e^{i\lambda x} \sum_j^0 f_j(x) + \sum_j^1 \left\{ \frac{\partial}{\partial z_1} \int_{x_*-1}^{x+z_1} f_j(s) J_0[\omega \sqrt{(x-s)^2 - z_1^2}] e^{i\lambda s} ds \right\} \\
= -e^{i\sigma} \frac{\partial}{\partial z_1} \int_{-1}^{x+z_1} (1+iks) J_0[\omega \sqrt{(x-s)^2 - z_1^2}] e^{i\lambda s} ds
\end{aligned}$$

$$\text{at } z_1 = -x_* \quad (V-14a)$$

and

$$\begin{aligned}
e^{i\lambda x} \sum_j^1 f_j(x) - \sum_j^0 \left\{ \frac{\partial}{\partial z} \int_{x_*-1}^{x-z} f_j(s) J_0[\omega \sqrt{(x-s)^2 - z^2}] e^{i\lambda s} ds \right\} \\
= \frac{\partial}{\partial z} \int_{-1}^{x-z} (1+iks) J_0[\omega \sqrt{(x-s)^2 - z^2}] e^{i\lambda s} ds \quad (V-14b)
\end{aligned}$$

at $z = x_*$ where

$$f_j(x) = 0 \text{ for } x \leq x_* - 1.$$

Performing the indicated differentiation yields

$$\begin{aligned} e^{i\lambda x} \sum_j \theta_j^0 f_j(x) + \sum_j \theta_j^1 \left\{ \int_{x_*-1}^{x-x_*} f_j(s) \frac{J_1[\omega \sqrt{(x-s)^2 - x_*^2}] \omega x_* e^{i\lambda s} ds}{\sqrt{(x-s)^2 - x_*^2}} \right. \\ \left. + f_j(x-x_*) e^{i\lambda(x-x_*)} \right\} \\ = e^{i\sigma} \left\{ - \int_{-1}^{x-x_*} (1+iks) \frac{\omega x_* J_1[\omega \sqrt{(x-s)^2 - x_*^2}] e^{i\lambda s} ds}{\sqrt{(x-s)^2 - x_*^2}} \right. \\ \left. - [1+ik(x-x_*)] e^{i\lambda(x-x_*)} \right\} \quad (V-15a) \end{aligned}$$

and

$$\begin{aligned} e^{i\lambda x} \sum_j \theta_j^1 f_j(x) + \sum_j \theta_j^0 \left\{ \int_{x_*-1}^{x-x_*} f_j(s) \frac{J_1[\omega \sqrt{(x-s)^2 - x_*^2}] \omega x_* e^{i\lambda s} ds}{\sqrt{(x-s)^2 - x_*^2}} \right. \\ \left. + f_j(x-x_*) e^{i\lambda(x-x_*)} \right\} \\ = - \int_{-1}^{x-x_*} (1+iks) \frac{\omega x_* J_1[\omega \sqrt{(x-s)^2 - x_*^2}] e^{i\lambda s} ds}{\sqrt{(x-s)^2 - x_*^2}} - [1+ik(x-x_*)] e^{i\lambda(x-x_*)} \quad (V-15b) \end{aligned}$$

Gorelov's formulation, equations [2.6, 2.7, 2.8, and 2.9] of [4], can be obtained directly from equations (V-15) by substituting

$$f_j(x) = \cos j\eta - \cos j\eta_*$$

$$v^0 = \sum_{j=0}^n \mu_j^0 \cos j\eta$$

$$v^1 = \sum_{j=0}^n \mu_j^1 \cos j\eta$$

where:

$$\eta = \cos(-x)$$

$$\eta_* = \cos(1-x_*)$$

In comparing the two systems care must be taken to note the differing symbols and coordinate systems. The corresponding quantities are:

Here	in [4]
θ_j^0, θ_j^1	$v_{0\sigma}, v_{i\sigma}$
μ_j^0, μ_j^1	$\theta_{0\sigma}, \theta_{i\sigma}$
$\eta = \cos^{-1}(-x)$	$\eta = \cos^{-1}(1-x)$
$\eta_* = \cos^{-1}(1-x_*)$	$\eta_* = \cos^{-1}(1-x_*)$
z, z_1	y, y_1
j	σ
σ	ψ

Here $-1 \leq x \leq 1$; in [4] $0 \leq x \leq 2$. This transformation accounts for the differing definitions of η . Making the substitutions results in the system

$$\sum_{j=0}^n \{ \theta_j^0 [\cos j\eta - (1-\delta_{01}) \cos j\eta_*] \\ + \theta_j^1 \int_{x_*-1}^{x-x_*} \frac{\partial}{\partial z_1} J_0 [\omega \sqrt{(x-s)^2 - z_1^2}] [\cos j\hat{\eta} - (1-\delta_{i0}) \cos j\eta_*] e^{i\lambda s} ds$$

$$+ \theta_j^1 [\cos j\bar{\eta} - (1-\delta_{i0}) \cos j\eta_*] e^{i\lambda(x-x_*)} \}$$

$$= - \sum_{j=0}^n \{ -\mu_j^1 \int_{-1}^{x-x_*} \frac{\partial}{\partial z_1} J_0 [\sqrt{(x-s)^2 - z_1^2}] \cos j\hat{\eta} e^{i\lambda s} ds \\ - \mu_j^1 \cos j\bar{\eta} e^{i\lambda(x-x_*)} \}$$

$$(V-16a)$$

$$\text{at } z_1 = -x_*$$

$$x > x_* - 1$$

and

$$\sum \{ \theta_j^1 [\cos j\eta - (1-\delta_{0j}) \cos j\eta_*]$$

$$+ \theta_j^0 \int_{x_*-1}^{x-x_*} \frac{\partial}{\partial z} J_0 [\omega \sqrt{(x-s)^2 - z^2}] [\cos j\eta - (1-\delta_{0j}) \cos j\eta_*] e^{i\lambda s} ds$$

$$+ \theta_j^0 [\cos j\bar{\eta} - (1-\delta_{0j}) \cos j\eta_*] e^{i\lambda(x-x_*)} \}$$

$$= \sum_{j=0}^n \{ \mu_j^0 \int_{-1}^{x-x_*} \frac{\partial}{\partial z} J_0 [\omega \sqrt{(x-s)^2 - z^2}] \cos j\hat{\eta} e^{i\lambda s} ds$$

$$- \mu_j^0 \cos j\bar{\eta} e^{i\lambda(x-x_*)} \} \quad (V-16b)$$

where:

$$\eta = \arccos(-x)$$

$$\eta_* = \arccos(1-x_*)$$

$$\bar{\eta} = \arccos(-s)$$

$$\hat{\eta} = \arccos(-x+x_*)$$

$$\delta_{0j} = \text{Dirac } \delta \text{ function} = \begin{matrix} 1 & \text{for } j = 0 \\ 0 & \text{for } j \neq 0 \end{matrix}$$

Given the change in coordinates and notation, this system is equivalent to that shown in [4].

This was the system programmed for computer solution. Because the function, $f_j(x)$, is unaffected by the differentiation with respect to y (or z) the exact form used need not be specified, so that the system shown in (V-15) may be programmed with f undetermined. A subroutine may be written to return the function desired and the remaining program left perfectly general. In the program developed with this thesis both the Gorelov functions shown above and the Legendre polynomials were employed. All the integrals may now be evaluated at $n+1$ points, x_i , on both blades in $x_{i-1} < x_i < 1$ and the resulting linear system solved for θ_j^0 and θ_j^1 , $j=1,2,\dots,n+1$.

The interference potentials may be constructed by taking

$$\psi^0(x,z) = \frac{-1}{m} \int_1^{x-z} [\theta_j^0 f_j(x)] J_0[\omega \sqrt{(x-s)^2 - z^2}] e^{i\lambda s} ds$$

$z > 0$ (V-17a)

$$\psi^1(x,z_1) = \frac{1}{m} \int_{-1}^{x+z_1} [\theta_j^1 f_j(x)] J_0[\omega \sqrt{(x-s)^2 - z_1^2}] e^{i\lambda s} ds$$

$z_1 < 0$ (V-17b)

where

$$f_j(x) = 0, \quad \text{for all } x \leq x_{*-1}$$

Once the potentials have been calculated as outlined above, the surface pressure may be calculated using the relationship

$$C_p = -2(\psi_x + ik\psi) \quad (V-18a)$$

$$= -2[\psi_x + i(k-\lambda)\psi]e^{-i\lambda x} \quad (V-18b)$$

Because all the plates are assumed to be in steady oscillation with uniform phase shift, σ , between neighboring plates, then

$$v^1(x) = v^0(x)e^{i\sigma}, \quad u^1(x) = u^0(x)e^{i\sigma}, \quad \psi(x,y) = -\psi(x,-y)$$

in this case

$$C_{\ell\alpha} = 2 \int_{-1}^1 [\psi_x(x,+0) + i(k-\lambda)\psi(x,+0)]e^{-i\lambda x} dx \quad (V-19)$$

where

$$\begin{aligned} \psi(x_1,+0) &= \frac{-1}{m} \int_{-1}^x [(1+iks)+u^0(s)]J_0[\omega(x-s)]e^{i\lambda(s)} ds \\ &+ \frac{e^{i\sigma}}{m} \int_{-1}^{x-x_*} [(1+iks)+u^0(s)]J_0[\omega\sqrt{(x-s)^2-(x_*)^2}]e^{i\lambda(s)} ds \end{aligned}$$

where $x_* = mp.$ (V-20)

Results from this approach, in the form of values of $C_{\ell\alpha}$ for $k = 0.1$, at various values of w are presented in the results for

approximations based both on Gorelov's formulation, and on the Legendre polynomials.

D. COLLOCATION SOLUTION OF THE POTENTIAL EQUATION EXPANDED FOR SMALL k ,

In order to provide a partially independent check of the results of the main program, the Gorelov function representation of the collocation solution was expanded for small k , and solved at two collocation points, $n = 2$. The resulting potentials, and partial derivatives with respect to x and y were then used to replace the corresponding numerical routines in the main program. The output resulting from the approximations were compared with the purely numerical results obtained from the computer program.

1. Solution For The Unknown Potential Coefficients

The basic system of linear equations used to determine the unknown coefficients is

$$\frac{1}{m} e^{i\lambda x} u^0 + \phi_{z_1}^1 + \psi_{z_1}^1 = 0, \quad z = 0, \quad z_1 = -x_* = -mp \quad (V-21a)$$

$$\frac{1}{m} e^{i\lambda x} u^1 + \phi_z^0 + \psi_z^0 = 0, \quad z_1 = 0, \quad z = x_* = mp \quad (V-21b)$$

where

$$0 = u^0(x) = u^1(x) \quad \text{when} \quad x \leq -1+x_*$$

otherwise

$$u^0 = \sum_{j=1}^n \theta_j^0 (\cos j\eta - \cos j\eta_*) + \theta_0^0$$

$$u^1 = \sum_{j=1}^n \theta_j^1 (\cos j\eta - \cos j\eta_*) + \theta_0^1$$

where

$$\eta = \arccos(-x)$$

$$\eta_* = \arccos(1 - x_*)$$

Thus, for $n = 2$, the system becomes

$$e^{i\lambda x} \{ \theta_0^0 + \theta_1^0 (\cos \eta - \cos \eta_*) \}$$

$$+ \frac{\partial}{\partial z_1} \int_{-1}^{x+z_1} v^1(s) J_0 [\omega \sqrt{(x-s)^2 - z_1^2}] e^{i\lambda s} ds$$

$$+ \frac{\partial}{\partial z_1} \int_{x_*-1}^{x+z_1} u^1(s) J_0 [\omega \sqrt{(x-s)^2 - z_1^2}] e^{i\lambda s} ds = 0$$

$$z_1 = -x_*$$

(V-22a)

along the reference blade and

$$e^{i\lambda x} \{ \theta_0^1 + \theta_1^1 (\cos \eta - \cos \eta_*) \} - \frac{\partial}{\partial z} \int_{-1}^{x-z} v^0(s) J_0 [\omega \sqrt{(x-s)^2 - z^2}] e^{i\lambda s} ds$$

$$- \frac{\partial}{\partial z} \int_{x_*-1}^{x-z} u^0(s) J_0 [\omega \sqrt{(x-s)^2 - z^2}] e^{i\lambda s} ds . \quad (V-22b)$$

$$z = x_*$$

along the adjacent blade

$$\text{where} \quad u^0(s) = \theta_0^0 + \theta_1^0 (\cos \eta - \cos \eta_*)$$

$$u^1(s) = \theta_0^1 + \theta_1^1 (\cos \eta - \cos \eta_*)$$

$$v^0(s) = 1 + iks$$

$$v^1(s) = (1 + iks) e^{i\sigma}$$

For k sufficiently small, this system may be further simplified by the following approximations

$$J_0 [\omega \sqrt{(x-s)^2 - z^2}] \approx 1 - O(\omega^2) \approx 1 \quad (V-23a)$$

$$J_0 [\omega \sqrt{(x-s)^2 - z^2}] \approx 1 - O(\omega^2) \approx 1 \quad (V-23b)$$

$$e^{i\lambda x} \approx 1 + i\lambda x - O(\lambda^2 x^2) \approx 1 + i\lambda x$$

$$e^{i\lambda s} \approx 1 + i\lambda s - O(\lambda^2 s^2) \approx 1 + i\lambda s$$

where $O(\omega^2)$ means "of the order of magnitude of ω^2 "

$$-1 \leq x \leq 1, \quad -1 \leq s \leq x - x_*$$

$$\lambda = k/m^2, \quad \omega^2 = \frac{k^2(1-m^2)}{m^4}$$

The interference source distributions may be replaced by

$$u^0(s) = \theta_0^0 + \theta_1^0(-s + x_* - 1)$$

$$u^1(s) = \theta_0^1 + \theta_1^1(-s + x_* - 1) .$$

If higher order terms are neglected, the result is a system linear in k and λ

$$\begin{aligned} (1+i\lambda x) [\theta_0^0 + \theta_1^0(-x-1+x_*)] + \frac{\partial e^{i\sigma}}{\partial z_1} \int_{-1}^{x+z_1} (1+iks)(1+i\lambda s) ds \\ + \frac{\partial}{\partial z_1} \int_{x_*-1}^{x+z_1} [\theta_0^1 + \theta_1^1(-s-1+x_*)](1+i\lambda s) ds = 0 \\ z_1 = -x_* \end{aligned} \quad (V-24a)$$

$$\begin{aligned} (1+i\lambda x) [\theta_0^1 + \theta_1^1(-x-1+x_*)] - \frac{\partial}{\partial z} \int_{-1}^{x-z} (1-iks)(1+i\lambda s) ds \\ - \frac{\partial}{\partial z} \int_{x_*-1}^{x-z} [\theta_0^0 + \theta_1^0(-s-1+x_*)](1+i\lambda s) ds = 0 \\ z = x_* \end{aligned} \quad (V-24b)$$

Product terms containing $(k\lambda) = \frac{k^2}{m^2}$ may be neglected as of higher order in k , yielding

$$(1+i\lambda x) [\theta_0^0 + \theta_1^0(-x-1+x_*)] + \frac{\partial e^{i\sigma}}{\partial z_1} \int_{-1}^{x+z_1} [1+i(\lambda+k)s] ds$$

$$+ \frac{\partial}{\partial z_1} \int_{x_*-1}^{x+z_1} \{ [\theta_0^1 + \theta_1^1(-s-1+x_*)] + i\lambda s [\theta_0^1 + \theta_1^1(-s-1+x_*)] \} ds = 0$$

$$z_1 = -x_* \quad (V-25a)$$

$$(1+i\lambda x) [\theta_0^1 + \theta_1^1(-x-1+x_*)] - \frac{\partial}{\partial z} \int_{-1}^{x-z} [1+i(\lambda+k)s] ds$$

$$- \frac{\partial}{\partial z} \int_{x_*-1}^{x-z} \{ [\theta_0^0 + \theta_1^0(-s-1+x_*)] + i\lambda s [\theta_0^0 + \theta_1^0(-s-1+x_*)] \} ds = 0$$

$$z = x_* \quad (V-25b)$$

Evaluating the indicated derivatives yields

$$(1+i\lambda x) [\theta_0^0 + \theta_1^0(-x-1+x_*)] + [1+i(k+\lambda)(x-x_*)] e^{i\sigma} + [\theta_0^1 + \theta_1^1(2x_*-x-1)] + i\lambda(x-x_*) [\theta_0^1 + \theta_1^1(2x_*-x-1)] = 0$$

$$(V-26a)$$

$$(1+i\lambda x) [\theta_0^1 + \theta_1^1 (-x-1+x_*)] + [1+i(\lambda+k)(x-x_*)]$$

$$+ \{ [\theta_0^0 + \theta_1^0 (2x_* - x - 1)] + i\lambda(x-x_*) [\theta_0^0 + \theta_1^0 (2x_* - x - 1)] \} = 0$$

(V-26b)

Thus:

$$\theta_0^0(1+i\lambda x) + \theta_1^0[(-x-1+x_*) + i\lambda x(-x-1+x_*)]$$

$$+ \theta_0^1[1 + i\lambda(x-x_*)] + \theta_1^1[(2x_* - x - 1) + i\lambda(x-x_*)(2x_* - x - 1)]$$

$$= -e^{i\sigma} [1 + i(k+\lambda)(x-x_*)] \quad (V-27a)$$

$$\theta_0^1[1 + i\lambda x] + \theta_1^1[(-x+1-x_*) + i\lambda x(-x+1-x_*)]$$

$$+ \theta_0^0[1+i\lambda x(x-x_*)] + \theta_1^0[(2x_* - x - 1) + i\lambda(x-x_*)(2x_* - x - 1)]$$

$$= -[1 + i(\lambda+k)(x-x_*)] \quad (V-27b)$$

This system may be solved at two points, x_1 and x_2 , for

θ_0^0 , θ_1^0 , θ_0^1 , and θ_1^1 .

2. Calculation Of The Potential

The potential is given by

$$\psi(x, y) = \Psi(x, z) e^{-i\lambda x} \quad (V-28)$$

where

$$\begin{aligned} \Psi(x, z) = & -\frac{1}{m} \int_{-1}^x [v^0(s) + u^0(s)] J_0[\omega \sqrt{(x-s)^2}] e^{i\lambda s} ds \\ & + \frac{1}{m} \int_{-1}^{x-x_*} [v^1(s) + u^1(s)] J_1[\omega \sqrt{(x-s)^2 - x_*^2}] e^{i\lambda s} ds. \end{aligned}$$

$$u^0(s) = u^1(s) = 0 \quad \text{for all } s \leq x_* - 1.$$

Thus

$$\begin{aligned} \Psi(x, z) = & -\frac{1}{m} \int_{-1}^x v^0(s) J_0[\omega \sqrt{(x-s)^2}] e^{i\lambda s} ds \\ & - \frac{1}{m} \int_{x_*-1}^x u^0(s) J_0[\omega \sqrt{(x-s)^2}] e^{i\lambda s} ds \\ & + \frac{1}{m} \int_{-1}^{x-x_*} v^1(s) J_0[\omega \sqrt{(x-s)^2 - x_*^2}] e^{i\lambda s} ds \\ & + \frac{1}{m} \int_{x_*-1}^{x-x_*} u^1(s) J_0[\omega \sqrt{(x-s)^2 - x_*^2}] e^{i\lambda s} ds \quad (V-29) \end{aligned}$$

Making the same small frequency approximations as in the previous section yields

$$\begin{aligned}
 \psi(x, z) = & -\frac{1}{m} \int_{-1}^x v^0(s) (1+i\lambda s) ds \\
 & -\frac{1}{m} \int_{x_*-1}^x u^0(s) (1+i\lambda s) ds \\
 & + \frac{1}{m} \int_{-1}^{x-x_*} v^1(s) (1+i\lambda s) ds \\
 & + \frac{1}{m} \int_{x_*-1}^{x-x_*} u^1(s) (1+i\lambda s) ds \quad (V-30)
 \end{aligned}$$

From the general formulation

$$\psi = \phi^1 + \phi^0 + \psi^1 + \psi^0 \quad (V-31)$$

Thus, along the reference blade

$$\begin{aligned}
 -m\phi^0(x, z=0) &= \int_{-1}^x v^0(s) (1+i\lambda s) ds = \int_{-1}^x (1+iks) (1+i\lambda s) ds \\
 &= \int_{-1}^x [1+i(k+\lambda)s] ds = \left[s + i(k+\lambda)\frac{s^2}{2} \right]_{-1}^x \quad (V-32) \\
 &= x + i\frac{(k+\lambda)}{2}x^2 + 1 - i\left(\frac{k+\lambda}{2}\right)
 \end{aligned}$$

$$\phi^0(x, z=0) = -\frac{1}{m}[(1+x) + i(\frac{k+\lambda}{2})(x^2-1)] \quad (V-33)$$

By inspection

$$\phi^1(x, z_1=-x_*) = \frac{e^{i\sigma}}{m}\{(1+x-x_*) + i(\frac{k+\lambda}{2})[(x-x_*)^2-1]\} \quad (V-34)$$

$$-m\psi^0(x, z=0) = \int_{x_*-1}^x u^0(s)(1+i\lambda s)ds \quad (V-35)$$

$$= \int_{x_*-1}^x [\theta_0^0 + \theta_1^0(-s+x_*-1)](1+i\lambda s)ds$$

$$= \int_{x_*-1}^x \theta_0^0(1+i\lambda s) + \theta_1^0(-s+x_*-1)(1+i\lambda s)ds$$

$$= \theta_0^0[(x+1-x_*) + \frac{i\lambda}{2}(x^2-x_*^2+2x_*-1)]$$

$$+ \int_{x_*-1}^x \theta_1^0[-x+x_*-1+i\lambda(-s^2+sx_*-s)]ds$$

$$= \theta_0^0[(x+1-x_*) + \frac{i\lambda}{2}(x^2-x_*^2+2x_*-1)]$$

$$+ \theta_1^0\{(-\frac{s^2}{2}+sx_*-s)+i\lambda(-\frac{s^3}{3}+\frac{s^2x_*}{2}-\frac{s^2}{2})\} \Big|_{x_*-1}^x \quad (V-36)$$

$$\begin{aligned}
&= \theta_0^0 \left[(x+1-x_*) + \frac{i\lambda}{2} (x^2 - x_*^2 + 2x_* - 1) \right] \\
&\quad + \theta_1^0 \left\{ \left(-\frac{x^2}{2} + x_* - x \right) - \left[-\frac{(x_*-1)^2}{2} + x_* (x_*-1) - (x_*-1) \right] \right. \\
&\quad \left. + i\lambda \left[\left(-\frac{x^3}{3} + \frac{x^2 x_*}{2} - \frac{x^2}{2} \right) + \frac{(x_*-1)^3}{3} - x_* \frac{(x_*-1)^2}{2} + \frac{(x_*-1)^2}{2} \right] \right\} \\
&= \theta_0^0 \left[(x+1-x_*) + \frac{i\lambda}{2} (x^2 - x_*^2 + 2x_* - 1) \right] \\
&\quad + \theta_1^0 \left\{ -\frac{x^2}{2} + x(x_*-1) + \frac{(x_*-1)^2}{2} - (x_*-1)^2 \right. \\
&\quad + i\lambda \left[-\frac{x^3}{3} + x^2 \frac{(x_*-1)}{2} + \frac{x_*^3}{3} - x_*^2 + x_* - \frac{1}{3} \right. \\
&\quad \left. \left. - \frac{x_*^3}{2} + x_*^2 - \frac{x_*}{2} + \frac{x_*^2}{2} - x_* + \frac{1}{2} \right] \right\} \\
&= \theta_0^0 \left[(x+1-x_*) + \frac{i\lambda}{2} (x^2 - x_*^2 + 2x_* - 1) \right] \\
&\quad + \theta_1^0 \left\{ -\frac{x^2}{2} + x(x_*-1) - \frac{(x_*-1)^2}{2} \right. \\
&\quad \left. + i\lambda \left[-\frac{x^3}{3} + \frac{x^2 (x_*-1)}{2} - \frac{x_*^3}{6} + \frac{x_*^2}{2} - \frac{x_*}{2} + \frac{1}{6} \right] \right\}
\end{aligned}$$

$$\begin{aligned}
\psi^0 = & -\frac{1}{m} \{ \theta_0^0 [(x+1-x_*) + \frac{i\lambda}{2}(x^2 + x_*^2 + 2x_*-1)] \\
& + \theta_1^0 \{ -\frac{x^2}{2} + x(x_*-1) - \frac{(x_*-1)^2}{2} \\
& + i\lambda [-\frac{x^3}{3} + \frac{x^2(x_*-1)}{2} - \frac{(x_*-1)^3}{6}] \} \} \quad (V-37)
\end{aligned}$$

$m\psi^1(x, z_1 = -x_*)$ may be evaluated by substituting $x-x_*$ for x in the expression for $m\psi^0$ and exchanging θ_0^1 and θ_1^1 for θ_0^0 and θ_1^0

$$\begin{aligned}
m\psi^1(x, z_1 = -x_*) = & \theta_0^1 \{ [(x-x_*)+1-x_*] + \frac{i\lambda}{2} [(x-x_*)^2 - x_*^2 + 2x_* - 1] \} \\
& + \theta_1^1 \{ -\frac{(x-x_*)^2}{2} + (x-x_*)(x_*-1) - \frac{(x_*-1)^2}{2} \\
& + i\lambda [-\frac{(x-x_*)^3}{3} + \frac{x^2(x_*-1)}{2} - \frac{(x_*-1)^3}{6}] \} \quad (V-38)
\end{aligned}$$

$$\begin{aligned}
= & \theta_0^1 (x+1-2x_*) + \frac{i\lambda}{2} [(x^2-2xx_*+x_*^2) - x_*^2 + 2x_* - 1] \} \\
& + \theta_1^1 \{ -\frac{(x^2-2xx_*+x_*^2)}{2} + (x - \frac{3}{2}x_* + \frac{1}{2})(x_*-1) \\
& + i\lambda [-\frac{x^3}{3} + x^2x_* - xx_*^2 + \frac{x_*^3}{3} + (x^2-2xx_*+x_*^2)\frac{(x_*-1)}{2} \\
& - \frac{(x_*-1)^3}{6}] \}
\end{aligned}$$

$$\begin{aligned}
&= \theta_0^1 \{ (x+1-2x_*) + \frac{i\lambda}{2} [x^2 - 2xx_* + 2x_* - 1] \\
&\quad + \theta_1^1 \{ [-\frac{x^2}{2} + xx_* - \frac{x_*^2}{2} + x(x-1) + \frac{1}{2}(1-3x_*)(x_*-1) \\
&\quad + i\lambda [-\frac{x^3}{3} + x^2 \frac{(3x_*-1)}{2} + x(-\frac{x_*^2-2xx_*^2-1)}{2} \\
&\quad \quad + \frac{x_*^3}{3} + \frac{x_*^3}{2} - \frac{1}{2}] \}
\end{aligned}$$

$$\psi^1(x, z_1 = -x_*)$$

$$\begin{aligned}
&= \frac{1}{m} \{ \theta_0^1 \{ (x+1-2x_*) + \frac{i\lambda}{2} [x^2 - 2xx_* + 2x_* - 1] \\
&\quad + \theta_1^1 \{ [-\frac{x^2}{2} + x(2x_*-1) - xx_*^2 + 2x_* - \frac{1}{2}] \\
&\quad + i\lambda [-\frac{x^3}{3} + x^2 \frac{(3x_*-1)}{2} + x \frac{(-3x_*^2-1)}{2} + 5\frac{x_*^3}{6} - \frac{1}{2}] \} \}
\end{aligned}$$

(V-39)

A comparison of the results for the full program and the approximation is given below for $k = 0.01$, $w = 0.05$, $\sigma = \rho$, $n = 2$, yielding $\omega^2 \approx 6.1 \times 10^{-3}$, $\lambda \approx 0.083$

	Full program	Approx
$x = .1285$		
ϕ	-5.363, .327i	-5.379, .549i
ϕ_x	-8.634, .361i	-8.66, .379i
$x = .5643$		
ϕ	-9.754, .809i	-9.804, 1.000i
ϕ_x	-12.234, .692i	-12.272, .7608i
C_{ℓ_α}	+31.658, -3.1032i	C_{ℓ_α} 31.7019, -3.2165i

VI. RESULTS

The collocation method was used to solve the partial differential equation resulting from the Gorelov approximation of transonic potential flow in an unstaggered cascade. The system was solved using both the spanning functions proposed by Gorelov in [4], resulting in the equations (V-16); and the Legendre polynomials, resulting in equations (V-15) with f_j replaced by the Legendre polynomial, P_j . The resulting values of $C_{l\alpha}$ for $k = .1$, $\tau = 1$, $\sigma = 1$ and seven collocation points on each blade are presented in figures VI-2, VI-3, and VI-4.

Figure VI -1 presents a diagram which is useful in commenting on the other results. This shows the location of the collocation points and first three interference reflections as a function of w expressed as a percentage of that portion of the chord subject to reflection. The collocation points are equally spaced throughout this interval, 12.5% from the leading edge of the interference zone, 12.5% between each pair of points and 12.5% from the blade trailing edge. The independent variable, w , is plotted vertically so that the dependent variable, percent of chord subject to interference, may be more conveniently visualized along the blade. (The curves are not precisely linear, but are very nearly so in the range shown.)

Figure VI -2 shows the $C_{l\alpha}$ calculated with $k = 0.0$ in comparison with the results obtained from Ackeret theory.

Agreement is good where there is no reflection and the portion of the blade subject to interference is affected by a constant interference potential, $w \geq 0.11$. Throughout the rest of the curve the results calculated here oscillate above and below the theoretical values. This appears to be due to the discrete nature of the approximation used in the collocation method. Rarely is the fraction of the chord subject to interference reflection equal to the fraction of the collocation points which feel it. Where the collocation point fraction lags, as near $w = 0.6$, the collocation results are lower than those due to Ackeret theory. When the collocation point fraction leads, as it does for $w \leq .04$ and briefly for $w \approx .08$, the the collocation results are higher than those due to Ackeret theory. The fault appears to be an intrinsic feature of the small number of points sampled. This results in a set of coefficients similar to those which would be obtained from a generalized Fourier series based on the integration of the Taylor series expansion about each point. This obviously cannot be a good approximation when both the function and its derivative are discontinuous at the reflections.

Figures VI -3 and VI -4 show the results of using Legendre polynomials and Gorelov's functions as spanning functions. The results for both formulations are identical. Gorelov's results are presented for comparison. Agreement is good for $w > 0.05$ except for an anomalous point, marked A .

It is believed that this anomaly is due to the location of the first reflection just ahead of the last collocation point (cf. "A" on Figure VI -1). This will yield a very small contribution from the reflection potential to the linear system from which the collocation points are determined. The resulting system will have a large dynamic range and may be ill-conditioned.

The discrepancy between these results, and those in [3] for $w < 0.5$ is still unexplained, as is the outlying value for $w = 0.5$.

The discontinuities in the imaginary results are believed to be due primarily to the reflection/collocation interaction explained above.

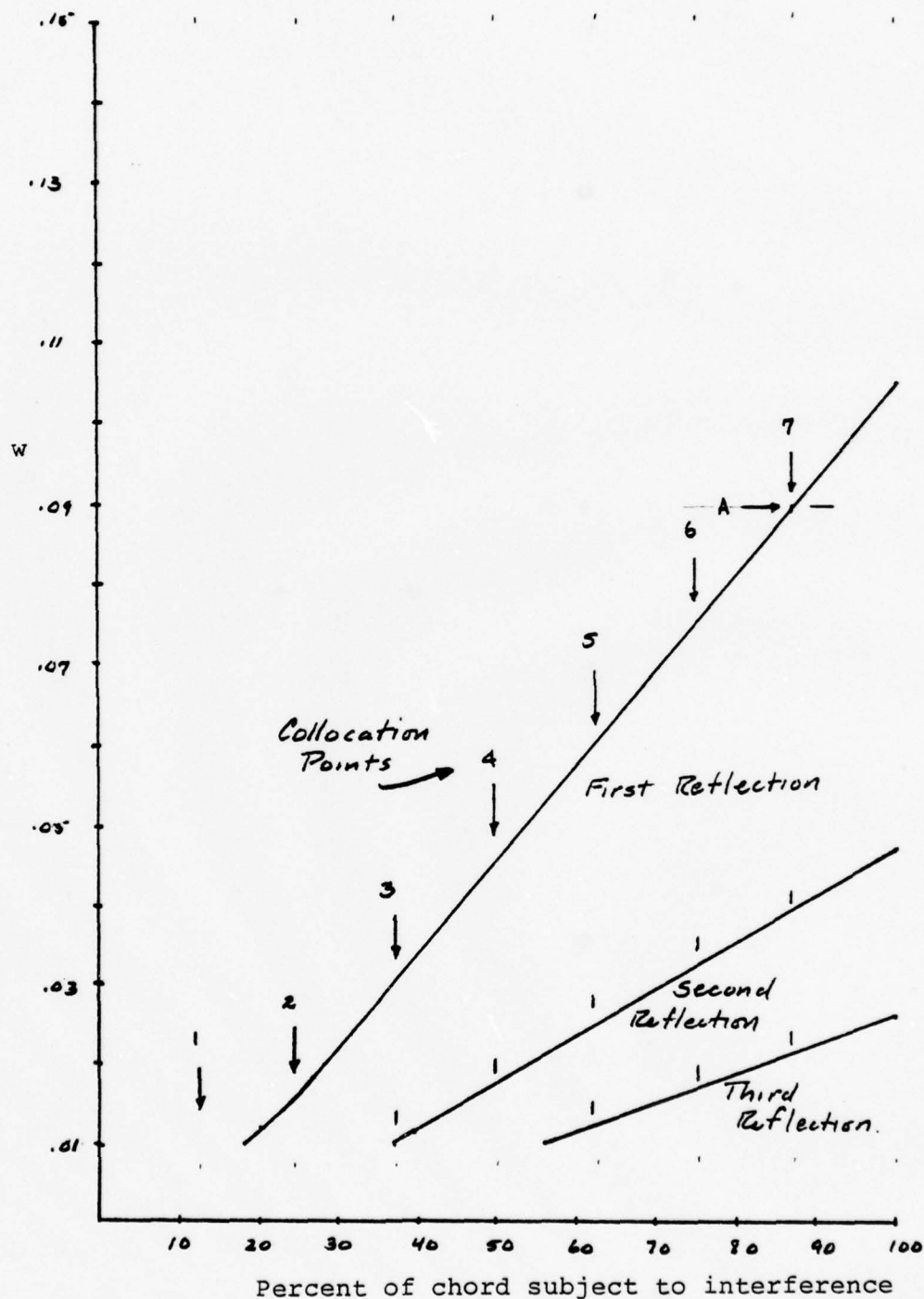


FIGURE VI-1. Location of Reflections and Collocation Points Shown as Percent of Chord Subject to Interference

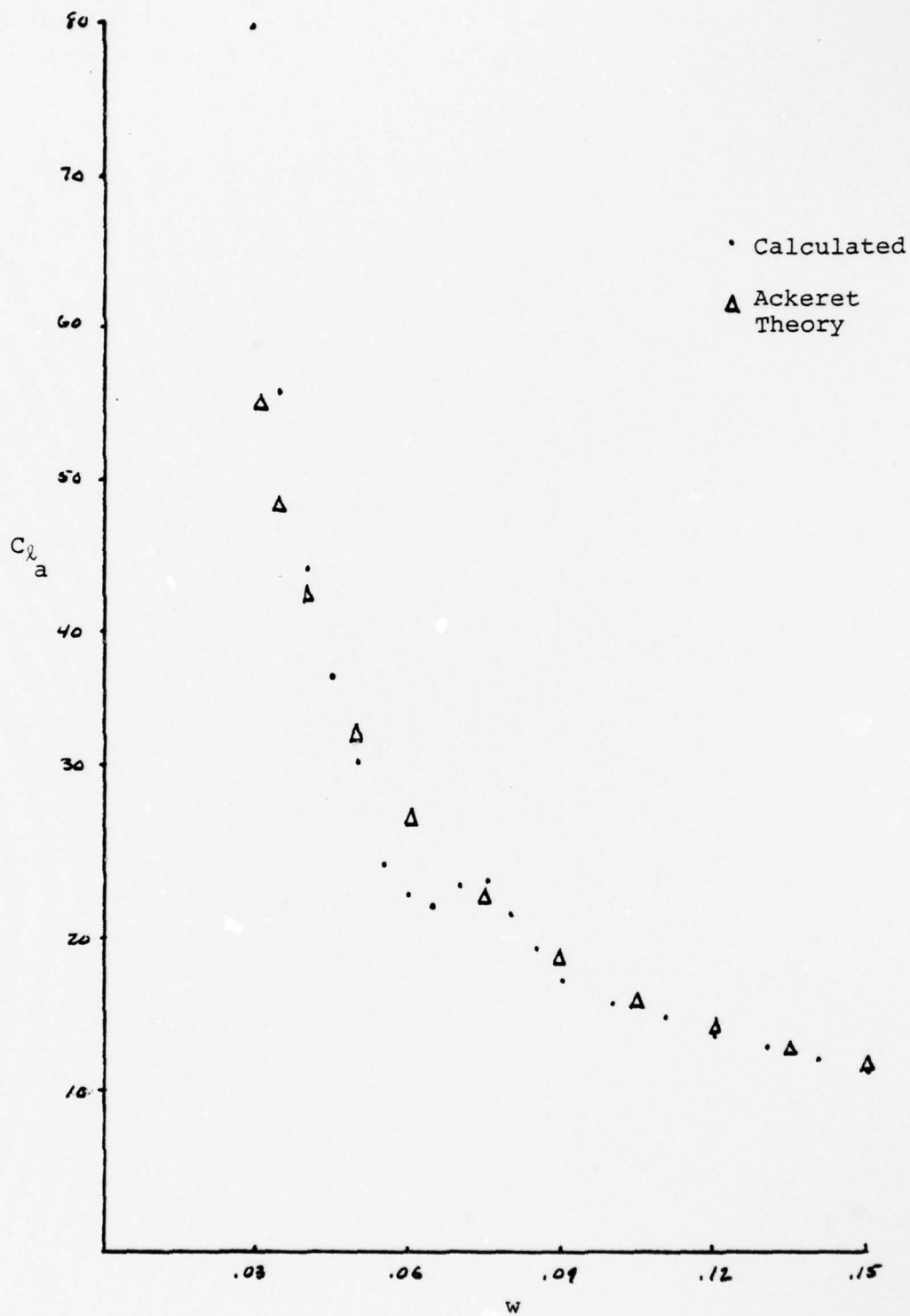


FIGURE VI-2. Comparison of C_{l_α} -vs- w to that Obtained from Ackeret Theory

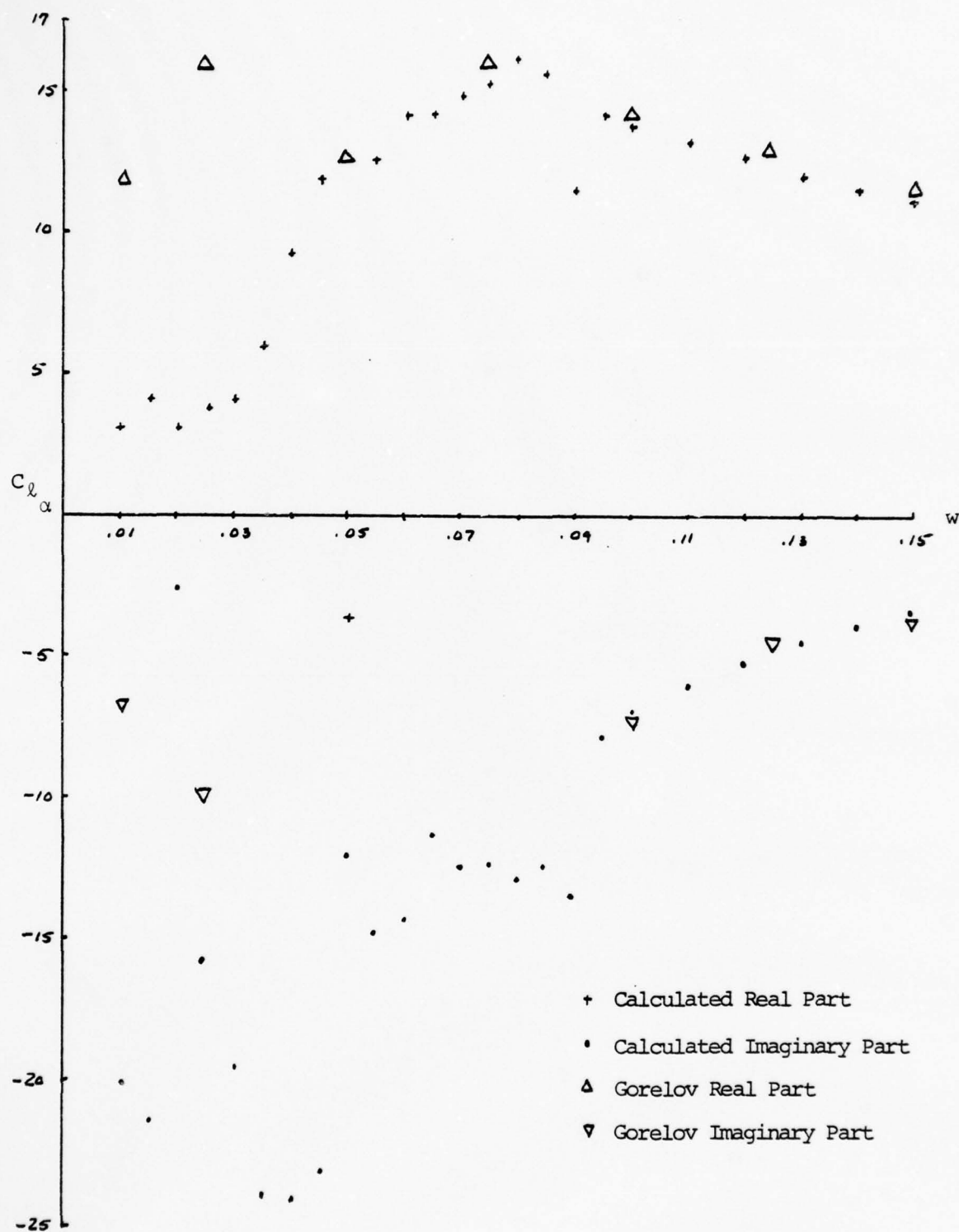


FIGURE VI-3. Plot of $C_{l\alpha}$ -vs- w , Legendre Polynomials
 $k = 0.1, \tau = 1.0, \sigma = \pi, n = 7$
 compared with Gorelov's results

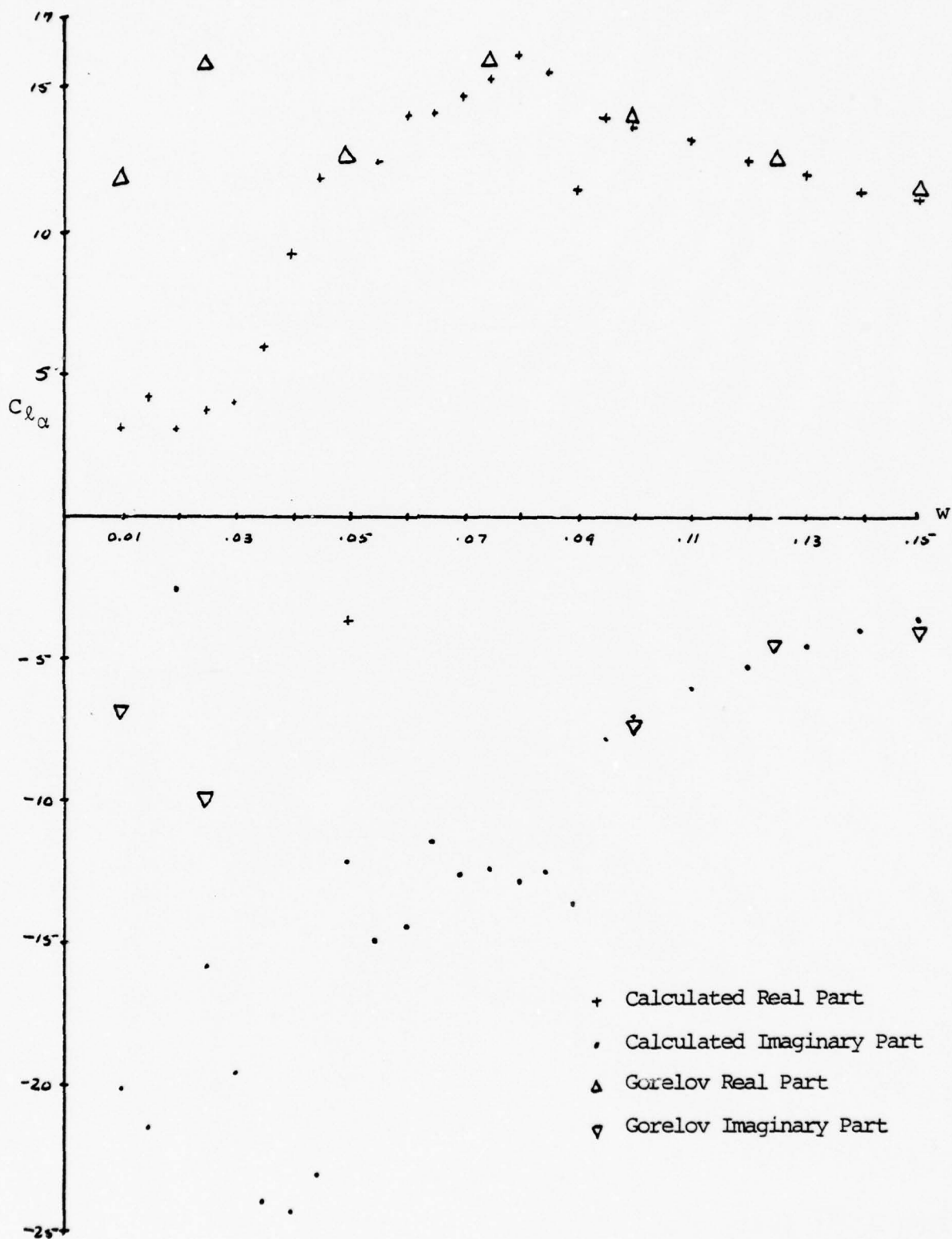


FIGURE VI-4. Plot of $C_{l\alpha}$ -vs- w , Gorelov Spanning Function
 $k = 0.1$, $\tau = 1.0$, $\sigma = \pi$, $n = 7$
 compared with Gorelov's results

VII . RECOMMENDATIONS

There are two recommendations to be made about the techniques used in the collocation method, and a new area in which it might be employed.

The program developed in the course of writing this thesis employs adaptive Simpson's integration to calculate the elements of a completely determined system. This system is solved to provide the coefficients of the spanning functions. Two improvements may be made:

1. The Simpson's integration scheme may be replaced by a Gaussian integrator. Experience has shown that several thousand function evaluations are required by the Simpson's integration routine when C_{t_α} is to be evaluated for small w . This entails large amounts of computer time and leads to increased accumulations of numerical error. Use of Gaussian integration would probably improve both of these characteristics with little loss of accuracy.

2. The present program treats a completely determined system of dimension $2n+1$ by $2n+1$, and then solves that system to find the collocation coefficients. This procedure has worked satisfactorily in this thesis, but may not work as well at higher frequencies where the final linear system of equations may be ill-conditioned. As an alternative, it is recommended that the boundary conditions be applied at more than n points, say m points, where m is twice or three times as many points, and that the least squares technique be used to determine the

the collocation coefficients which give the minimum square error over-all. This may be thought of as "sampling more data" in order to get more information about the unknown function. The present program could be easily modified in this regard by replacing the spanning function matrix, Q1ZINT, by a new matrix of the form

$$Q1ZINT' = X^T X$$

where X is the new m by n+1 ($m > n+1$) matrix, and replacing the present right-hand-side vector, Q1COF with

$$Q1COF' = X^T Y$$

where Y is the new m by 1 right-hand-side vector. An alternative would be to employ a prepackaged statistical linear regression routine after either modifying the routine to accept complex data, or transforming the present system into a larger system of real numbers only.

The new area in which the collocation method might be employed is the calculation of the potential flow about a staggered cascade. The method could be employed to calculate both the potential in the channel and above the upper blade. The program presented has been designed to enable the

calculation of flow within the channel of a staggered cascade. Unfortunately, there was not enough time to extend the study to this case.

APPENDIX A
PROGRAM DESCRIPTION

This section describes the computer program used to calculate the interference solution to the Gorelov linearization for unsteady transonic flow in a channel. The program written in IBM Fortran IV with the basic structure outlined by Stevens [5]. The basic points are:

1. Organization of the program into small subroutines, each of which performs a specific task.
2. Transmission data to and from subroutines via a formal parameter argument list. No common statements are used.

The end objective is code which is both easy to modify and maintain.

Each subroutine is designed with optional diagnostic printing of its input and output. This is controlled by the parameter IPT. The diagnostic output is printed (only) if $IPT > 0$. Each routine accepts IPT, sets $IOT = IPT - 1$, and then passes IOT as the print parameter to routines it calls. By this method, diagnostic output can be "cascaded" to any desired level. Large initial values of IPT should be avoided because of the spectacular amount of output which can be generated by the double integrals within Q1DCOF.

1. Main Program; including subroutines READ and ABSA.

The basic structure is given in Figure A-1. MAIN calls READ to read input data and then ABSA to calculate the

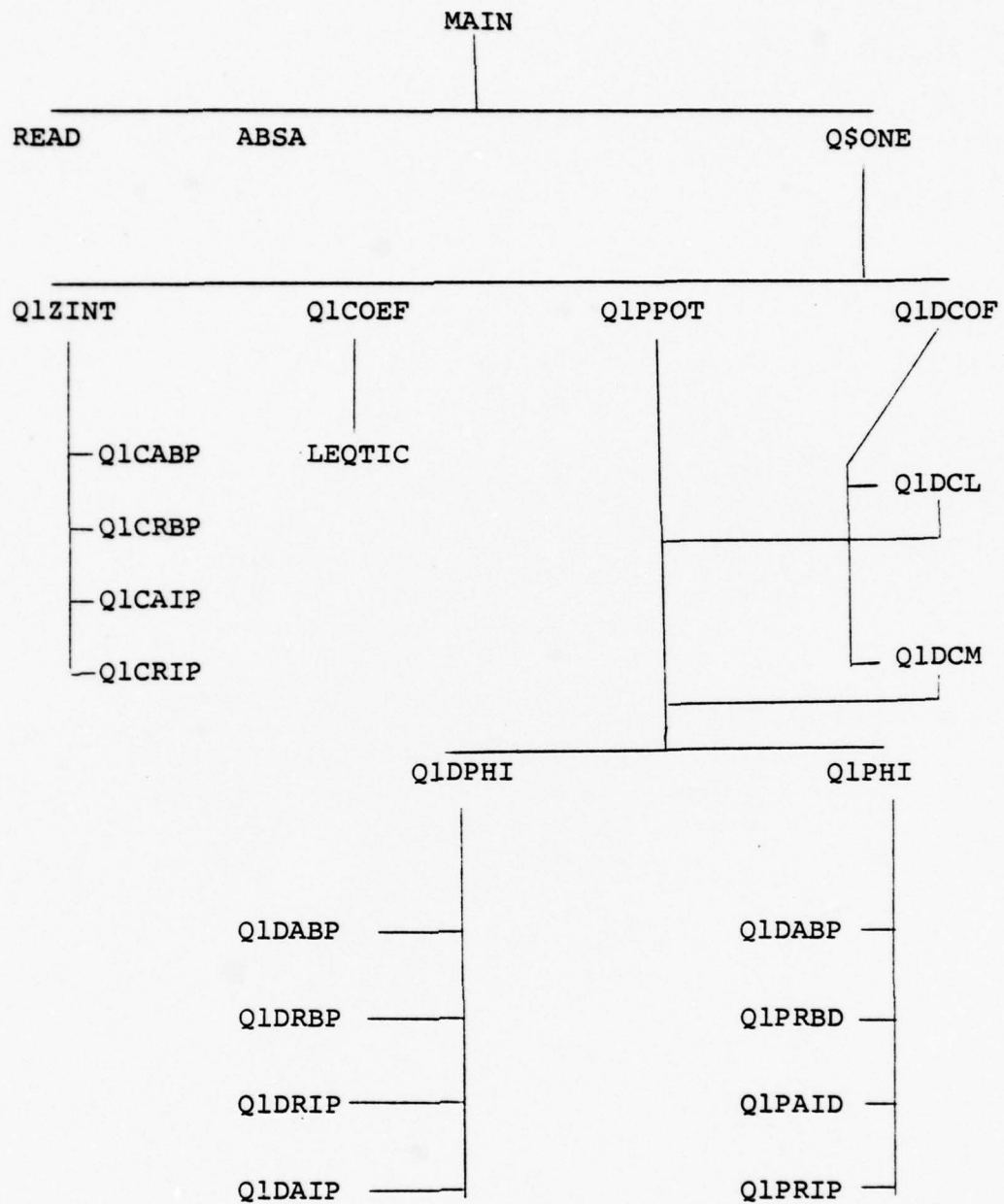


Figure A-1. Program Hierarchy

collocation points. The version of ABSA shown evenly spaces the collocation points across that portion of the blade subject to interference. ABSA may be easily replaced if different point spacing is desired, or if additional points are to be added for an overdetermined system and least squares approximation.

2. Q\$ONE This subroutine controls the actual potential calculation. It performs no calculation itself, but calls subordinate subroutines where the calculations are actually performed. The calling hierarchy is shown in Figure A-1.

3. Q1ZINT This subroutine calculates the linear equation system arising from the boundary conditions. Hierarchy is shown in Figure A-2.

The matrix output is carried through Q1INT. Q1ZINT calls the following subprograms

- a. Q1CRBP returns the value of ϕ_z^0
- b. Q1CABP returns the value of $\phi_{z_1}^1$
- c. Q1CRIP returns the values of ψ_z^0

$$\frac{\partial}{\partial z} \int_{-1+x_*}^{x-z} f_j(s) J_0[\omega \sqrt{(x-s)^2 - z^2}] ds$$

where f_i is one of the set of i elementary functions, $j=1,n$

- d. Q1CAIP returns the values of $\psi_{z_1}^1$

Q1ZINT

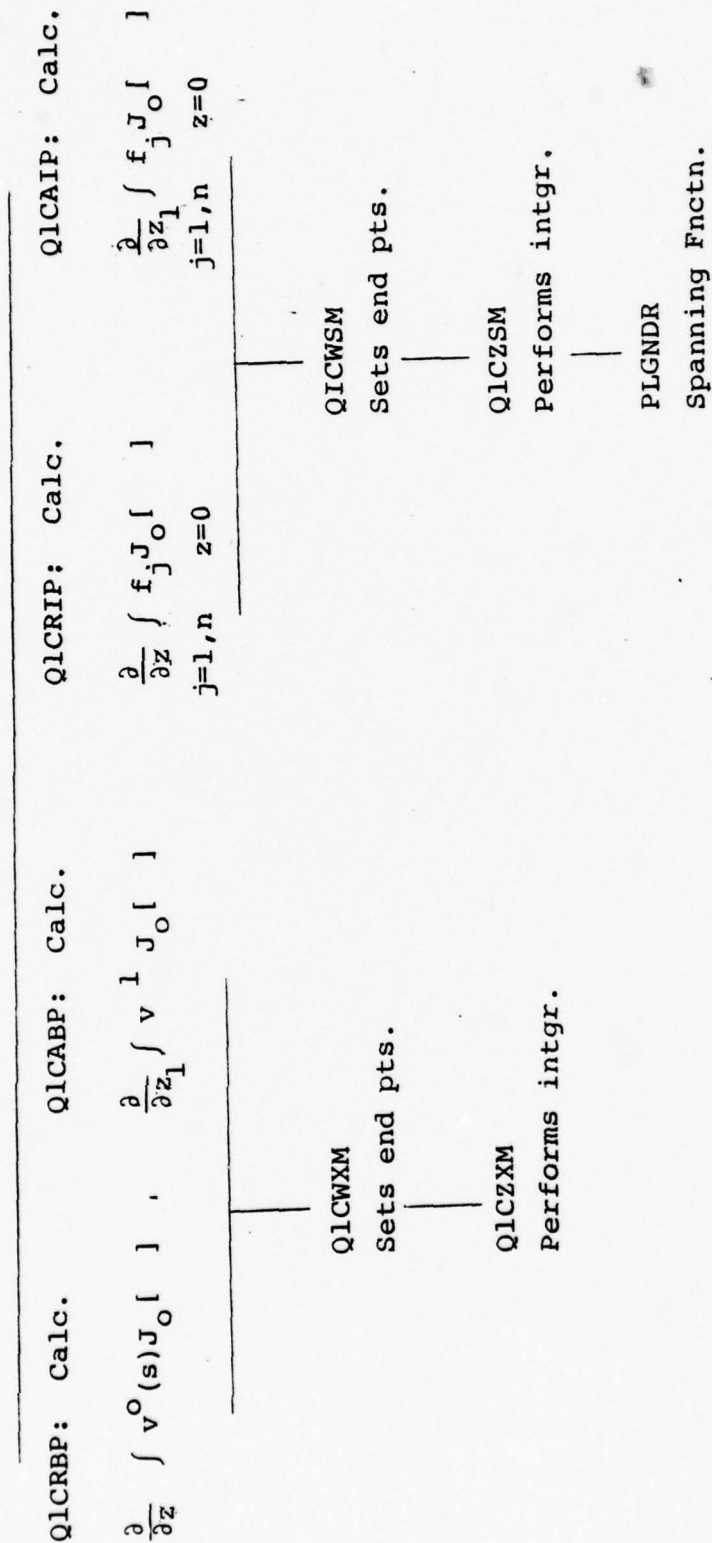


Figure A-2

$$\frac{\partial}{\partial z_1} \int_{-1+\text{OFFSET}+x_*}^{x+z_1} f_j(s) J_0[\omega \sqrt{(x-s)^2 - y_1^2}] ds$$

where f_j is one of the set of elementary functions. OFFSET is a parameter included to facilitate program conversion to a staggered cascade

e. PLGNDR is the subprogram which returns $f_j(x)$, the elementary spanning function. No other routine contains explicit reference to the spanning function. This facilitates easy replacement of the spanning functions should this be desired.

Q1CRBP and Q1CABP in turn call Q1CWXM and Q1CZXM. Q1CWXM computes end-points and then calls Q1CZXM, a complex integration routine based on SIMP by Shampine and Allen [6]. Q1CAIP and Q1CRIP call Q1CWSM and Q1CZSM to perform the integration. Q1ZINT passes the constant matrix to Q1COEF in the array Q1INT and the right-hand-side vector in the array Q1COF.

4. Q1COEF This subroutine employs the IMSL routine LEQT1C to solve the linear system received from Q1ZINT. The resulting coefficients are Q1ABCF for the adjacent blade and Q1RBCF for the reference blade. LEQT2C, the high precision complex IMSL routine may be directly substituted for LEQT1C. Q1COEF may be rewritten to employ the generalized inverse required for least squares approximation

$$A = (x^T x)^{-1} x^T y \text{ where } x = Q1INT$$

$$y = Q1COF$$

after first performing the multiplication necessary to replace Q1INT and Q1COF with the proper matrix products in the call to LEQTIC.

5. Q1PPOT This subroutine calculates the potential, ϕ , and ϕ_x , at each collocation point along the reference blade but only if Q1PPOT receives a value of IPT > 0, requiring IPT \geq 2 on input to the main program. If IPT \leq 0, then the subroutine is exited before any calculations are performed. This subroutine is most useful for debugging Q1ZINT and Q1COEF. Q1PPOT calls Q1PABP, Q1PRBP, Q1PRIP, Q1PAIP, Q1DABP, Q1DRBP, Q1DRIP, and Q1DAIP, all of which will be described in the next section.

6. Q1DCOF This subroutine calculates the dimensionless coefficients of lift and moment; C_{l_α} , C_{m_α} . Its internal hierarchy is shown in Figure A-3.

- a. Q1DCL calculates the nondimensional complex coefficient C_{l_α}
- b. Q1DCM calculates the nondimensional complex coefficient C_{m_α} .
- c. Q1PRBP and Q1PABP calculate the potentials due to the reference and adjacent blades respectively. Q1PWXM and Q1PZXM are called to perform the actual integration.

- d. Q1PRIP and Q1PAIP calculate the interference potentials along the reference and adjacent blades. Q1PZSM is called to perform the integration.
- e. Q1DRBP, Q1DABP, Q1DRIP, and Q1DAIP correspond exactly to subroutines above except that the value returned is the partial derivative of the potential with respect to X. Q1DWXM, Q1DZSM, and Q1DXSM perform the co-reponding integrals.

6. Program Listing The program listing shown below incorporates the Legendre functions as spanning functions. Listings for a subroutine employing Gorelov's spanning function and the linear approximation program follow.

QLDR0005
QLDR0010
QLDR0015
QLDR0020
QLDR0025
QLDR0030
QLDR0035
QLDR0040
QLDR0045
QLDR0050
QLDR0055
QLDR0060
QLDR0066C7C
QLDR0075
QLDR0080
QLDR0085
QLDR0090
QLDR0095
QLDR0100
QLDR0105
QLDR0110
QLDR0115
QLDR0120
QLDR0125
QLDR0130
QLDR0135
QLDR0140
QLDR0145
QLDR0150
QLDR0155
QLDR0160
QLDR0165
QLDR0170
QLDR0175
QLDR0180
QLDR0185
QLDR0190
QLDR0195
QLDR0200
QLDR0205
QLDR0210
QLDR0215
QLDR0220
QLDR0225
QLDR0230
QLDR0235
QLDR0240

CLDR C245
 CLDR C255
 CLDR C260
 CLDR C265
 CLDR C270
 CLDR C275
 CLDR C280
 CLDR C285
 CLDR C290
 CLDR C295
 CLDR C300
 CLDR C305
 CLDR C310
 CLDR C315
 CLDR C320
 CLDR C325
 CLDR C330
 CLDR C335
 CLDR C340
 CLDR C345
 CLDR C350
 CLDR C355
 CLDR C360
 CLDR C365
 CLDR C370
 CLDR C375
 CLDR C380
 CLDR C385
 CLDR C390
 CLDR C395
 CLDR C400
 CLDR C405
 CLDR C410
 CLDR C415
 CLDR C420
 CLDR C425
 CLDR C430
 CLDR C435
 CLDR C440
 CLDR C445
 CLDR C450
 CLDR C455
 CLDR C460
 CLDR C465
 CLDR C470
 CLDR C475
 CLDR C480

```

SUBROUTINE ABSA (N,OFFSET,X,DR,CW)
  IMPLICIT REAL*8 (A-H,O,P,R-Y), COMPLEX*16(Q,Z)
  DIMENSION X(13)
  XINT = (2.000-DR)/DFLOAT(N+1)
  XL = DR - 1.000 + 1.00D-8
  CC 10 I = 1,N
  X(I) = XL + (XINT*DFLOAT(I))
  IF(X(I).EQ.0.000) X(I) = .1E-14
  CONTINUE
10 RETURN
END
SUBROUTINE $ONE (CK,CR,CW,RHO,OFFSET,SIGMA,N,NF,IPT,X)
  IMPLICIT REAL*8(A-H,O,P,R-Y), COMPLEX*16(Z,C)
  DIMENSION Q1COF(26), Q1INT(26,26), Q1RBP(13), Q1ABP(13)
  DIMENSION Q1RBCF(13), Q1ABCF(13)
  DIMENSION X(13)
  IF(IPT.GT.0)WRITE(6,990)DK,CR,RHO,OFFSET,SIGMA,N,NF,IPT
  IPT = IPT - 1
  FORMAT(1,0)
1. 5X, 'DK', 13X, 'RHC', 12X, 'OFFSET', 9X, 'SIGMA', 10X,
2. 'N', 3X, 'NF', 2X, 'IPT', 2X, 'RHO', 12X, '2X', 13X
3. 5X, 'Q1ZINT', 12X, '2X', 13X
  CALL Q1ZINT(DK,DR,CW,RHO,OFFSET,SIGMA,N,X,IPT,C1CCF,C1INT)
  CALL Q1COEF(Q1COF,Q1INT,N,IPT,Q1AECF,Q1RBCF)
  CALL Q1PCT(DK,DR,CW,RHO,OFFSET,SIGMA,N,X,C1AECF,C1RBCF,IPT)
  C T = 0
  CALL Q1CCCF(CK,DR,CW,RHO,OFFSET,SIGMA,N,Q1AECF,Q1RBCF,IPT)
  RETURN
END
C
C
SUBROUTINE Q1ZINT(DK,DR,CW,RHO,OFFSET,SIGMA,N,X,
  IPT,Q1COF,Q1INT)
  IMPLICIT REAL*8(A-H,O,P,R-Y), COMPLEX*16(Z,Q)
  DIMENSION Q1COF(26), Q1INT(26,26), Q1INTRP(13), Q1INTAP(13)
  DIMENSION CN X(13)
  FORMAT(1,0,10X, 'Q1ZINT ENTERED WITH:', 12X, 'OFFSET', 9X, 'SIGMA', 10X,
  1. 'N', 2X, 'IPT', 2X, 'RHO', 12X, '2X', 13X
  2. 'N', 2X, 'IPT', 2X, 'RHO', 12X, '2X', 13X
  3. 1CX, 5(E12.5, 3X, 1, 12, 2X, 13)
  CCNST = CDEXP(DCMPLX(0.300, SIGMA))
  IF(IPT.GT.0) WRITE(6,990)DK,DR,RHO,OFFSET,SIGMA,N,IPT
  IPT = IPT - 1
  GAMMA = 1.40
  CLAMDA = 1.000/(GAMMA+1.000)*CW
  CC 90 I = 1,N
  IN = I + N
990

```



```

XSTN = X(I)
XQ = XSTN - DR
CEXP = CDEXP(DCMPLX(0.0D0, DLAMCA*XSTN))
CALL Q1CAIP(DK, DR, DW, RHO, OFFSET, SIGMA, XSTN, QINTAP, N, IOT)
CALL Q1CRIP(DK, DR, DW, RHO, OFFSET, XSTN, QINTRP, N, ICT)
DC 20 J = 1, N
JN = N + J
JI = J - 1
QIINT(I, J) = PLGNDR(XSTN, DR, JI) * CEXP
QIINT(I, J) = QINTRP(I, J)
QIINT(I, JN) = QINTAP(J)
QIINT(I, JN) = PLGNDR(XSTN - OFFSET, CR, JI) * CEXP
CCCONTINUE
QICOF(I, N) = -QICREF(DK, DR, DW, RFO, XSTN, IOT)
QICOF(I) = -QICABP(DK, DR, DW, RHO, OFFSET, SIGMA, XSTN, IOT)
CCCONTINUE
RETURN
END
CCOMPLEX FUNCTION QICRBP*16(DK, DR, [L, RFO, XSTN, IPT)
IAPLICIT REAL * 8 (A-F, C, P, R-Y), COMPLEX * 16 (Q, Z)
DIMENSION QINP(2)
IF (IPT.GT.0) WRITE(6,990) DK, DR, DW, RHO, XSTN, IPT
FORMAT(0, '10X, QICREP ENTERED WITH:', //,
1, '10X, CK, 11X, DR, 11X, DW, 11X, RFO, 10X, XSTN, 5X, IPT,
2, '/', XSTN, LE, RFO-1.0D0) GOTO 20
IF(XSTN.GT.2.0D0) GOTO 20
ICT = IPT - 1
QCK = DCMPLX(0.0D0, DK)
CALL Q1CWM(DK, DR, DW, RHO, XSTN, QINP, IOT)
QICRBP = -QCK*QINP(2) - QINP(1)
IF(IPT.LE.0) RETURN
GOTO 30
QICRBP = DCMPLX(0.0D0, 0.0C0)
IF(IPT.LE.0) RETURN
WRITE(6,995) QICRBP
FORMAT(0, '10X, QICREP = ', E14.7, ', ', E14.7)
RETURN
END
CCOMPLEX FUNCTION Q1CABP*16(DK, DR, DW, RHO, OFFSET, SIGMA, XSTN, IPT)
IAPLICIT REAL * 8 (A-F, O, P, R-Y), COMPLEX * 16 (C, Z)
DIMENSION QINP(2)
IF(IPT.GT.0) WRITE(6,990) DK, CR, DW, RHO, OFFSET, SIGMA, XSTN, IPT
FORMAT(0, '10X, Q1CABP ENTERED WITH:', //,
1, '10X, OK, 11X, CR, 11X, DW, 11X, RHO, 10X, XSTN, 5X, IPT, /,
2, 'SIGMA, 8X, XSTN, 10X, 7(E12.5, ', ', 13)
XSTN = XSTN - OFFSET
IF(XSTN.LE.RHO-1.0C0) GOTO 20

```

QLCR0485
 QLCR0490
 QLCR0495
 QLCR0500
 QLCR0505
 QLCR0510
 QLCR0515
 QLCR0520
 QLCR0525
 QLCR0530
 QLCR0535
 QLCR0540
 QLCR0545
 QLCR0550
 QLCR0555
 QLCR0560
 QLCR0565
 QLCR0570
 QLCR0575
 QLCR0580
 QLCR0585
 QLCR0590
 QLCR0595
 QLCR0600
 QLCR0605
 QLCR0610
 QLCR0615
 QLCR0620
 QLCR0625
 QLCR0630
 QLCR0635
 QLCR0640
 QLCR0645
 QLCR0650
 QLCR0655
 QLCR0660
 QLCR0665
 QLCR0670
 QLCR0675
 QLCR0680
 QLCR0685
 QLCR0690
 QLCR0695
 QLCR0700
 QLCR0705
 QLCR0710
 QLCR0715
 QLCR0720

84

QDR12105
QDR12115
QDR12125
QDR12135
QDR12145
QDR12155
QDR12165
QDR12175
QDR12185
QDR12195
QDR12205
QDR12215
QDR12225
QDR12235
QDR12245
QDR12255
QDR12265
QDR12275
QDR12285
QDR12295
QDR12305
QDR12315
QDR12325
QDR12335
QDR12345
QDR12355
QDR12365
QDR12375
QDR12385
QDR12395
QDR12405
QDR12415
QDR12425
QDR12435
QDR12445

```

FIT(LVL) = FV(3)
F2T(LVL) = FV(4)
F3T(LVL) = FV(5)
DAT(LVL) = DX
AREST(LVL) = ARESTL
ARESTT(LVL) = ARESTR
GESTT(LVL) = GESTL
GESTR(LVL) = QESTR
EPS(LVL) = EPS/1.4
EPS(LVL) = EPS
FV(5) = FV(3)
FV(3) = FV(2)
GC TO 1
ERROR = QERRR + QDIFF/15.C
IF(LORR(LVL).EQ.0) GO TO 4
CSUM = QPSUM(LVL) + CSUM
LVL = LVL - 1
IF(LVL.GT.1) GO TO 3
QANS = CSUM - (B*J1) * CDEXP(QEXF*B)
IF(IPT.EQ.0) GO TO 11
IF(IER.EQ.129) GO TO 11
IF(FLAG.EQ.1) RETURN
WRITE(6,990) DK,DR,DW,RHO,XSTN,A,B,J,IPT
FCRMAT(1,1,15X,13,2X,13,5X,E14.7,/,E14.7,/,
1,1,15X,13,2X,13,5X,E14.7,/,E14.7)
2,1,15X,13,2X,13,5X,E14.7,/,E14.7)
RETURN
QPSUM(LVL) = QSUM
LCRR(LVL) = 1
ALPHA = ALPHA + DA
DA = DAT(LVL)
FV(1) = FIT(LVL)
FV(3) = F2T(LVL)
FV(5) = F3T(LVL)
AREST = ARESTT(LVL)
EPS = EPST(LVL)
GC TO 1 2
IFLAG = 2 2
GC TO 2 3
IFLAG = 2 3
GC TO 2
END
SUBROUTINE QCAIP(CK,CR,DW,RTO,CFST,SGMA,XSTN,CCAIP,N,IPT)
IMPLICIT REAL*8 (A-H,C,P,R-Y), COMPLEX * 16 (C,Z)

```



```

DC 40 J = 1,N
WRITE (6,556) J,QCRIP(J)
FCRMAT(1,1,26X,12,3X,E14.7,1,1,E14.7)
40 CC CONTINUE
RETURN
ENC
SUBROUTINE QICWSM (DK,DR,DW,RHO, XSTN,N,QINF,IPT)
N IS THE MAXIMUM DEGREE OF THE TERM U IN THE INTEGRALS
IMPLICIT REAL * 8 (A-H,O,P,R-Y), COMPLEX * 16 (Z,C)
DIMENSION QINP(13)
IF (IPT.GT.0) WRITE (6,990) DK,DR,FHC,XSTN,N,IPT
FCRMAT(1,1,10X,QICWSM ENTERED WITH ARGUMENTS:1,/,
1,1,10X,DK,16X,DR,16X,RHO,15X,XSTN,14X,N,6X,IPT,/,
2,1,10X,4(E13.6,5X),12,5X,13)
IOT = IPT - 1
B = XSTN - RHO - 1.0D-10
A = DR - 1.0D0
DC 30 J = 1,N
CALL QICZSM(DK,DR,DW,RHO,XSTN,A,B,J,QANS,IOT)
QINF(J) = QANS
CC CONTINUE
IF (IPT.LE.0) RETURN
WRITE (6,995)
FCRMAT(1,1,10X,QICWSM RESULTS:1,/,
1,1,10X,J,6X,QINP(1),)
DC 1001 I = 1,N
FCRMAT(1,1,10X,13,5X,E12.6,1,1,E12.6)
556 WRITE (6,996) I, QINP(I)
1001 RETURN
END
SUBROUTINE QICZSM (DK,DR,DW,RHO, XSTN,A,B,J,QANS,IPT)
IMPLICIT REAL*8 (A-H,O,P,R-Y), COMPLEX*16 (F,Q,Z)
DIMENSION FV(5),LCRR(30),F1T(30),F2T(30),F3T(30),F4T(30),
1 ARESTT(30),QESTT(30),EPST(30),QPSUM(30)
F (X) = PLGNCR(X,DR,J1)*CDEXP(QEXP(X))
1 *(CMEGA*RHO/(LSQRT((XSTN-X)*(XSTN-X)-YY)))*
2 MBSJ1(OMEGA*DSQRT((XSTN-X)*(XSTN-X)-YY),IER))
GAMMA = 1.4D0
CM2 = (GAMMA + 1.0D0) * DW
CMEGA = CSQRT(DK*DK*(1.0D0-DM2)/(CM2*CM2))
YV = RHC * RHO
CLAMDA = DK/DM2
QEXP = DCMPLX(0.0D0,[LA MDA])
J1 = J-1
ACC = 1.0D-6
U = 9.0D-13

```

```

QLCR1685
QLCR1690
QLCR1695
QLDR1700
QLDR1705
QLDR1710
QLDR1715
QLDR1720
QLDR1725
QLDR1730
QLDR1735
QLDR1740
QLDR1745
QLDR1750
QLDR1755
QLDR1760
QLDR1765
QLDR1770
QLDR1775
QLDR1780
QLDR1785
QLDR1790
QLDR1795
QLDR1800
QLDR1805
QLDR1810
QLDR1815
QLDR1820
QLDR1825
QLDR1830
QLDR1835
QLDR1840
QLDR1845
QLDR1850
QLDR1855
QLDR1860
QLDR1865
QLDR1870
QLDR1875
QLDR1880
QLDR1885
QLDR1890
QLDR1895
QLDR1900
QLDR1905
QLDR1910
QLDR1915
QLDR1920

```

```

IF (IFT.GT.0) WRITE(6,990)DK,DR,CW,RHO,XSTN,A,E,J,IPT
990 FCRMAT(0,15X,'Q1CZSN ENTERED WITH ARGUMENTS:',/,10X,'RHO',12X,'XSTN',11X,
1,15X,'DK',13X,'CR',13X,'CW',10X,'RHO',12X,'XSTN',11X,
2,14X,'B',14X,'J',12X,'IPT',/,
3,15X,'7(E14.7,.,),12,2X,13),
FPCURU = 4.*U
IFLAG = 1
EPS = ACC
CERROR = DCMPLX(0.0D0,0.0D0)
LVL = 1
LCRR(LVL) = 1
CFSUM(LVL) = 0.0
ALPHA = A
DA = B - A
AREA = C.0
AREST = 0.0
FV(1) = F(ALPHA)
FV(3) = F(ALPHA + 0.5*DA)
FV(5) = F(ALPHA + DA)
KCUNT = 3
WT = DA/6.0
QEST = WT*(FV(1) + 4.0*FV(3) + FV(5))
DX = 0.5*DA
FV(2) = F(ALPHA + 0.5*DX)
FV(4) = F(ALPHA + 1.5*DX)
KCUNT = KCUNT + 2
WT = DX/6.0
QESTL = WT*(FV(1) + 4.0*FV(2) + FV(3))
QESTR = WT*(FV(3) + 4.0*FV(4) + FV(5))
QSUM = QESTL + QESTR
ARESTL = WT*(CDABS(FV(1)) + CDAES(4.0*FV(2)) + CDABS(FV(3)))
ARESTR = WT*(CDABS(FV(3)) + CDABS(4.0*FV(4)) + CDABS(FV(5)))
AREA = ARESTL + ARESTR - AREST
QDIFF = QEST - QSUM
IF(CDABS(QDIFF).LE.EPS*CDABS(AREA))GO TO 2
IF(CDABS(CX).LE.EFOURU*CDABS(ALPHA))GO TO 5
IF(LVL.GE.3)GO TO 5
IF(KCUNT.GE.2000)GC TC 6
LVL = LVL + 1
LCRR(LVL) = 0
F1T(LVL) = FV(3)
F2T(LVL) = FV(4)
F3T(LVL) = FV(5)
DA = DX
DAT(LVL) = CX
ARESTL = ARESTL + ARESTR
AFESTT(LVL) = ARESTL
QEST = QESTL

```

1

Q1DR1525
 Q1DR1530
 Q1DR1535
 Q1DR1540
 Q1DR1545
 Q1DR1550
 Q1DR1555
 Q1DR1560
 Q1DR1565
 Q1DR1570
 Q1DR1575
 Q1DR1580
 Q1DR1585
 Q1DR1590
 Q1DR1595
 Q1DR1600
 Q1DR2005
 Q1DR2010
 Q1DR2015
 Q1DR2020
 Q1DR2025
 Q1DR2030
 Q1DR2035
 Q1DR2040
 Q1DR2045
 Q1DR2050
 Q1DR2055
 Q1DR2060
 Q1DR2065
 Q1DR2070
 Q1DR2075
 Q1DR2080
 Q1DR2085
 Q1DR2090
 Q1DR2095
 Q1DR2100
 Q1DR2105
 Q1DR2110
 Q1DR2115
 Q1DR2120
 Q1DR2125
 Q1DR2130
 Q1DR2135
 Q1DR2140
 Q1DR2145
 Q1DR2150
 Q1DR2155
 Q1DR2160

```

100 CESTT(LVL) = QESTR
101 EPS = EPS/1.4
102 EFFST(LVL) = EPS
103 FV(5) = FV(3)
104 FV(3) = FV(2)
105 GO TO 1
106 QERROR = QERROR + CDIFF/15.0
107 IF(LCRR(LVL).EQ.0) GO TO 4
108 QSUM = QPSUM(LVL) + QSUM
109 LVL = LVL - 1
110 IF(LVL.GT.1) GO TO 3
111 CANS = FLGDCR(B,CR,J1) * CDEXP(QEXP*B) - QSUM
112 IF(IPT.GT.C) GO TO 11
113 IF(IER.EQ.129) GO TO 11
114 IF(IER.EQ.1) RETURN
115 WRITE(6,990) DK,DR,DW,RHC,XSTN,A,E,J,IPT
116 FCRMAT(1,15X,RESULTS,QANS,IFLAG,IER,QERR)
117 1. ,15X,IFLAG,2X,IER,5X,CERRCF,/,
118 2. ,15X,13,2X,13,5X,E14.7,.,.,E14.7)
119 RETURN
120 QPSUM(LVL) = QSUM
121 LCRR(LVL) = 1
122 ALPHA = ALPHA + DA
123 DA = DAT(LVL)
124 FV(1) = F1T(LVL)
125 FV(3) = F2T(LVL)
126 FV(5) = F3T(LVL)
127 AREST = AREST(LVL)
128 CEST = QEST(LVL)
129 EPS = EPST(LVL)
130 GO TO 1
131 IF LAG = 2
132 IF LAG = 3
133 IF LAG = 2
134 GO TO 6
135 REAL FUNCTION PLGNDR*(X,CR,N)
136 IMPLICIT REAL*8(A-T,O-Z)
137 IF(N.EQ.0) GO TO 100
138 X2 = X*X
139 GO TO (101,102,103,104,105,106,107,108,109,110,111,112), N
140 PLGNDR = 1.000
141 RETURN
142 PLGNDR = X
143 RETURN
144 PLGNDR = ((3.00C)*X2-1.00C)/2.00C
145 RETURN

```

CLDR2165
 CLDR2170
 CLDR2175
 CLDR2180
 CLDR2185
 CLDR2190
 CLDR2195
 CLDR2200
 CLDR2205
 CLDR2210
 CLDR2215
 CLDR2220
 CLDR2225
 CLDR2230
 CLDR2235
 CLDR2240
 CLDR2245
 CLDR2250
 CLDR2255
 CLDR2260
 CLDR2265
 CLDR2270
 CLDR2275
 CLDR2280
 CLDR2285
 CLDR2290
 CLDR2295
 CLDR2300
 CLDR2305
 CLDR2310
 CLDR2315
 CLDR2320
 CLDR2325
 CLDR2330
 CLDR2335
 CLDR2340
 CLDR2345
 CLDR2350
 CLDR2355
 CLDR2360
 CLDR2365
 CLDR2370
 CLDR2375
 CLDR2380
 CLDR2385
 CLDR2390
 CLDR2395
 CLDR2400

```

103 PLGNDR = ((5.000*X2 - 3.000)*X/2.000
RETURN
104 PLGNDR = ((3.501*X2 - 3.001)*X2 + 3.000)/8.000
RETURN
105 PLGNCR = ((6.301*X2 - 7.001) *X2 +1.501)*X/8.000
RETURN
106 PLGNDR=(((231.000*X2-315.000)*X2+105.000)*X2-5.000)/16.000
RETURN
107 PLGNDR=(((429.000*X2-693.000)*X2+315.000)*X2-35.000)
1 *X/16.000
RETURN
108 PLGNDR=(((6435.000*X2-1202.000)*X2+6930.000)*X2
1 -1260.000)*X2+35.000)/128.000
RETURN
109 PLGNDR=(((12155.000*X2-25740.000)*X2+18018.000)
1 *X2-4620.000)*X2+315.000)/128.000
RETURN
110 PLGNDR=(((146185.000*X2-109395.000)*X2+50000.000)
1 *X2-30030.000)*X2-3465.000)*X2-63.000)/256.000
RETURN
111 PLGNDR=(((188175.000*X2-230945.000)*X2+218750.000)*X2
1 -90090.000)*X2+15015.000)*X2-693.000)*X2/256.000
RETURN
112 PLGNDR=(((1676039.000*X2-1939535.000)*X2+2078505)*X2
1 -1021020.000)*X2+225225.000)*X2-18018.000)*X2-231.000)/1024.000
RETURN
ENCL
SLBRCTINE C1COEF( Q1COF, Q1INT, N, IPT, Q1ABCF, Q1RBCF)
IMPLICIT REAL*8(A-H,O,P,R-Y), COMPLEX*16(Z,C)
DIMENSION C1ABCF(26), C1INT(26,26), ZHA(300)
IF (IPT.GE.C) WRITE (6,90)
IE = 1
N2 = 1
N2 = 2*N
IF (IPT.LE.0) GO TO 5
WRITE (6,98) N,N2
58 FCRMAT(1,5X,'Q1COEF ENTERED WITH ',I2,' DEG PWR SERIES (',I2,
1 ', SQUARE MATRIX)',
DC 2 I = 1,N2
WRITE(6,92) I, I, Q1COF(I)
92 FCRMAT(1,10X,'Q1COEF EQUATION SYSTEM, ROW ',I2,/,
1 ',',10X,'C1COF( ',I2,') = ',E14.7, ',',E14.7)
DC 2 J = 1,N
J2=J+N
WRITE(6,91) I,J,Q1INT(I,J),I,J2,C1INT(I,J2)
91 FCRMAT(1,15X,2('Q1INT( ',I2, ', ',I2, ') = ',E14.7, ', ',E14.7,10X))
2 CONTINUE

```

```

Q1DR24010
Q1DR24115
Q1DR24220
Q1DR24325
Q1DR24430
Q1DR24535
Q1DR24640
Q1DR24745
Q1DR24850
Q1DR24955
Q1DR25060
Q1DR25165
Q1DR25270
Q1DR25375
Q1DR25480
Q1DR25585
Q1DR25690
Q1DR25795
Q1DR25800
Q1DR25905
Q1DR26010
Q1DR26115
Q1DR26220
Q1DR26325
Q1DR26430
Q1DR26535
Q1DR26640

```



```

5  IA = 26
   IJCB=0
   CALL LECTIC(QIINT,N2,IA,QICCF,M,IE,IJOB,ZWA,IER)
   IF(IER.EQ.0) GO TO 30
   IF(IER.EQ.129) GO TO 10
   WRITE(6,93)
93  FORMAT('0',10X,'QICCEF - ITERATIVE IMPROVEMENT FAILED, MATRIX TOO
      1  ILL-CONDITIONED. USE RESULTS WITH CAUTION.')
   CALL WRITE(6,95)
1C  WFORMAT('0',10X,'QICCEF - MATRIX ALGORITHMICALLY SINGULAR. CCEFFIC
95  IENTS SET TO ZERO.')
   ZERO = DCMPLX(0.0DC,C.0D0)
   DO 20 I = 1,N
   I2 = I+N
   QICCOF(I) = ZERO
   QICCOF(I2) = ZERO
   QICCONTINUE
2C  DO 30 I = 1,N
3C  IN = I + N
   QIABCF(I) = QICCOF(IN)
   QIRBCF(I) = QICCOF(I)
   QICCONTINUE
35  IF(IPT.LE.0) RETURN
   IF(IPT.GE.0) WRITE(6,94)
   DO 40 I = 1,N
   IM1 = I-1
   WRITE(6,99) I, IM1, QIRBCF(I), QIABCF(I)
40  CCNTINUE
55  FCRMAT('0',5X,I2,10X,I2,10X,2(E14.7,5X,E14.7,10X))
5C  FCRMAT('0',10X,'SUBROUTINE QICCEF - COMPLEX POWER SERIES COEFFIC
54  FCRMAT('0',5X,'INDEX',7X,'DEG FCLY',4X,'REFERENCE BLADE TERMS QIC
3  FCRF(INDEX),7X,'ADJACENT BLADE TERMS QICOF(2*INDEX),')
   RETURN
   ENCL
   FUNCTION IQIIFC(I)
   .....
   IQIIFC IS A FUNCTION SUBROUTINE DESIGNED TO PRODUCE FACTORIALS
   .....
   IF(I.EQ.0) GO TO 11
   IF(I.EQ.1) GO TO 11
   IQIIFC=I
   IF I=1
   DO 10 J=1,I
10  IQIIFC=IQIIFC*(I-J)

```

CC


```

QDRIPP = QIDRIP(DK,DR,DW,RHO,OFFSET,XSTN,QIRBCF,N,IOT)
QCAIPO = QIDAIPO(DK,DR,DW,RZERO,OFFSET,SIGMA,XSTN,QIABCF,N,IOT)
QDAIPP = QIDAIPO(DK,DR,DW,RHO,OFFSET,SIGMA,XSTN,QIABCF,N,IOT)
QCR = QDRP + QDAF + QCRIPQ + QCAIPP
QDA = QDRP + QDAO + QCRIPQ + QDAIFC
QCPHI(I) = QIDCFI(DK,DR,DW,RHO,OFFSET,SIGMA,N,QIABCF,CIRBCF,
1 XSTN,IOT)
WRITE(6,55) I,XSTN
WRITE(6,91) QCR,QCRO,QDAP,QDRIPQ,CCAIPP
WRITE(6,92) CCA,QCRP,QDAO,QCRIPP,CCAIPQ
QCRRI(I) = QCR
QCAA(I) = CCA
FCRMMAT(1,0,X-STATIC NUMBER,12,XSTN = F6.4,/,
1 BL TCTAL D(POT)/DX,14X,REF BL D(POT)/DX,8X,
2 ADJ BL D(FCT)/DX,8X,REF BL INT D(POT)/DX,5X,
3 ADJ BL INT D(POT)/DX,1)
10 CONTINUE
95 WRITE(6,94) XSTN,CFHI(I),QRR(I),QAA(I)
FCRMMAT(1,1,10X,SUMMARY LISTING,/,
954 1,0,10X,XSTN,7X,SINGLE BLADE TOTAL POTENTIAL,6X,
1 REF BLADE POTENTIAL,15X,ADJ BLADE POTENTIAL,1)
DC 20 I = 1,N
XSTN = X(I)
WRITE(6,94) XSTN,CFHI(I),QRR(I),QAA(I)
FCRMMAT(1,1,10X,F6.4,3(5X,E14.7,/,E14.7))
20 CONTINUE
96 WRITE(6,96)
FCRMMAT(1,0,10X,XSTN,7X,SINGLE ELADE TOTAL D(PCT)/DX,6X,
956 1 REF ELADE D(POT)/DX,15X,ADJ ELADE D(POT)/DX,1)
DC 50 I = 1,N
XSTN = X(I)
WRITE(6,94) XSTN,CCPHI(I),QCRRI(I),CDAI(I)
50 CONTINUE
RETURN
END
SUBROUTINE QIDCOF(DK,DR,DW,RHC,CFFSET,SIGMA,N,QIABCF,QIRBCF,IPT)
IMPLICIT REAL*8 (A-H,O,P,R-Y)
COMPLEX*16 (Q,Z)
DIMENSION QIABCF(12),QIRBCF(13)
IF(IPT.GT.0)WRITE(6,990) DK,DR,C,RHO,OFFSET,SIGMA,N,IPT
990 FCFMMAT(1,1,10X,QIDCOF - CALCULATION OF COMPLEX DIMENSIONLESS AERO-
1 DYNAMIC COEFFICIENTS,/,
3 0,10X,13X,DR,13X,DW,13X,RHO,12X,OFFSET,5X,SIGMA,
4 10X,N,3X,IPT,/,
56(E12.5,3X),12,2X,13)
ICCT = IPT - 3
995 QCCCL = QIDCCL(DK,DF,DW,RHO,OFFSET,SIGMA,N,QIABCF,CIRBCF,IOT)
QCCM = QIDCCM(DK,DR,DW,RHC,CFFSET,SIGMA,N,QIABCF,QIRBCF,IOT)
CANMA = 1.4CO

```

QDR31125
 QDR31130
 QDR31135
 QDR31140
 QDR31145
 QDR31150
 QDR31155
 QDR31160
 QDR31165
 QDR31170
 QDR31175
 QDR31180
 QDR31185
 QDR31190
 QDR31195
 QDR3200
 QDR3205
 QDR3210
 QDR3215
 QDR3220
 QDR3225
 QDR3230
 QDR3235
 QDR3240
 QDR3245
 QDR3250
 QDR3255
 QDR3260
 QDR3265
 QDR3270
 QDR3275
 QDR3280
 QDR3285
 QDR3290
 QDR3295
 QDR3300
 QDR3305
 QDR3310
 QDR3315
 QDR3320
 QDR3325
 QDR3330
 QDR3335
 QDR3340
 QDR3345
 QDR3350
 QDR3355
 QDR3360

```

TAL = (2.0D0*DSQRT((GAMMA + 1.0D0)*DW))/DR
WRITE (6,90) DK,TAU,DW,N,SIGMA,CCCL,CCDM
C FCRMAT(0,5X,DK = ,F6.3,, F7.4,, DW = ,F6.3,, N = ,
1 I2,, SIGMA = ,F6.3,, CL = ,F9.4,, F9.4,, CM = ,
2 F9.4,, F9.4)
RETURN
END
COMPLEX FUNCTION QIDCL*16(DK,DR,DW,RHO,OFFSET,SIGMA,N,QIABCF,CIRBCF,
1 I,IMPLICTION QIABCF(13), QIRBCF(13), COMPLEX*16(F,Q,Z)
DIMENSION FV(5),FIT(60),FET(60),QEST(60),CPSUM(60)
DIMENSION DAT(60),AREST(60),EPST(60)
DIMENSION LCRR(60)
F(X) = (C1DEFI(DK,DR,CW,RHO,OFFSET,SIGMA,N,QIABCF,X,ICT)
1 -QALPHA*CLPHI(CK,DR,CW,RHO,OFFSET,SIGMA,N,QIABCF,X,ICT)
0 3 IF (IPT.GT.C) WRITE (6,990) CK,CR,DW,RHO,CFFSET,SIGMA,N,IPT
C FCRMAT(0,10X,DK,13X,DR,13X,CK,13X,RHO,12X,OFFSET,9X,SIGMA,
4 10X,5,3X,12,2X,13)
5 (E12.5,3X),12,2X,13)
A = 1.0D0
ECT = 1.0D0
ICT = 1.0D0
GAMMA = 1.4D0
CLAMDA = CK/((GAMMA + 1.0D0) * DW)
QALPHA = DCMLPX(0.0D0,CLAMDA-DK)
U ACC = 9.0E-13
AEFCURU = 1.0D-5
IFLAG = 1
EPS = ACC
QERROR = 1
QVLR(LVL) = 1
QPSUM(LVL) = 0.0
ALPHA = A
DA = B - A
AREAT = 0.0
AREST = 0.0
FV(1) = F(ALPHA)
FV(3) = F(ALPHA) + 0.5*DA
FV(5) = F(ALPHA + DA)
KCLNT = 3
MT = DA/6.0
CEST = MT*(FV(1) + 4.0*FV(3) + FV(5))
DX = 0.5*DA
FV(2) = F(ALPHA + 0.5*DX)

```

QIDCL = 65
 QIRBCF = 70
 QIABCF = 75
 QIDCL = 80
 QIRBCF = 85
 QIABCF = 90
 QIDCL = 95
 QIRBCF = 100
 QIABCF = 105
 QIDCL = 110
 QIRBCF = 115
 QIABCF = 120
 QIDCL = 125
 QIRBCF = 130
 QIABCF = 135
 QIDCL = 140
 QIRBCF = 145
 QIABCF = 150
 QIDCL = 155
 QIRBCF = 160
 QIABCF = 165
 QIDCL = 170
 QIRBCF = 175
 QIABCF = 180
 QIDCL = 185
 QIRBCF = 190
 QIABCF = 195
 QIDCL = 200
 QIRBCF = 205
 QIABCF = 210
 QIDCL = 215
 QIRBCF = 220
 QIABCF = 225
 QIDCL = 230
 QIRBCF = 235
 QIABCF = 240
 QIDCL = 245
 QIRBCF = 250
 QIABCF = 255
 QIDCL = 260
 QIRBCF = 265
 QIABCF = 270
 QIDCL = 275
 QIRBCF = 280
 QIABCF = 285
 QIDCL = 290
 QIRBCF = 295
 QIABCF = 300

AD-A063 083

NAVAL POSTGRADUATE SCHOOL MONTEREY CALIF
THEORETICAL ANALYSIS OF TRANSONIC FLOW PAST UNSTAGGERED OSCILLA--ETC(U)
SEP 78 P C OLSEN

F/G 20/4

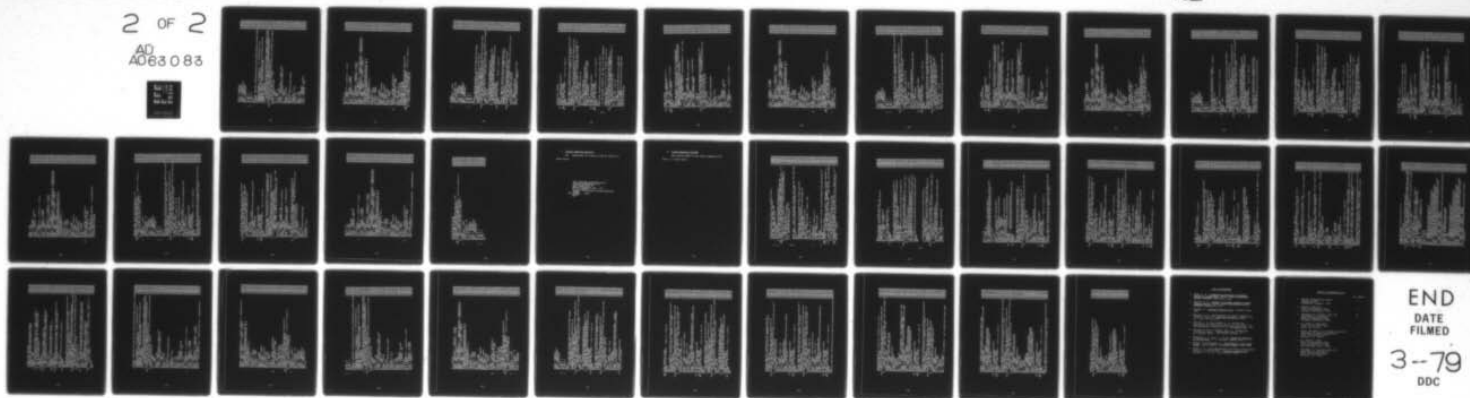
UNCLASSIFIED

2 OF 2

AD
A063 083

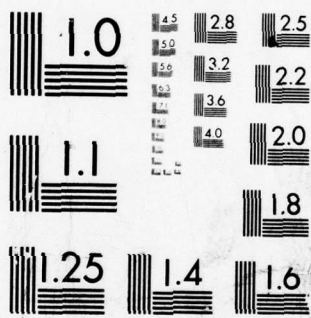


NL



END
DATE
FILMED

3--79
DDC



MICROCOPY RESOLUTION TEST CHART
NATIONAL BUREAU OF STANDARDS-1963-A


```

1      KCUNT = 3
        WT = DA/6.0 + FV(1) + 4.0*FV(3) + FV(5)
        CEST = 0.5*DA
        FV(2) = F(ALPHA + C.5*DX)
        FV(4) = F(ALPHA + 1.5*DX)
        KOUNT == KCUNT + 2
        DT = DX/6.0
        QESTL = WT*(FV(1) + 4.0*FV(2) + FV(3))
        CESTR = WT*(FV(3) + 4.0*FV(4) + FV(5))
        CSUM = CESTL + CESTR
        ARESTL = WT*(CDABS(FV(1))) + CDABS(4.0*FV(2)) + CDABS(FV(3)))
        ARESTR = WT*(CDABS(FV(3))) + CDABS(4.0*FV(4)) + CDABS(FV(5)))
        AREA = AREA + ((ARESTL + ARESTR) - AREST)
        CDIFF = QUEST - QSUM
        IF(CDABS(QDIFF).LE.EPS*DABS(AREA)){GO TO 2}
        IF(CDABS(DX).LE.EFCURU*DABS(ALPHA)){GO TO 3}
        IF(LVL.GE.60){GO TO 5}
        IF(KOUNT.GE.4000){GC TO 6}
        LVL = LVL + 1
        LCRR(LVL) = 0
        FLT(LVL) = FV(3)
        FZT(LVL) = FV(4)
        FBT(LVL) = FV(5)
        DAT = DX
        DAT(LVL) = CX
        ARESTL = ARESTR
        ARESTT(LVL) = ARESTR
        CESTT(LVL) = QUEST
        EPS = EPS/1.4
        EFST(LVL) = EPS
        FV(5) = FV(2)
        FV(3) = FV(1)
        GC TO 1
        QERRCR = QERROR + QCIF/15.0
        IF(LORR(LVL).EQ.0){GO TO 4}
        CSUM = QPSUM(LVL) + QSUM
        LVL = LVL - 1
        IF(LVL.GT.1){GO TO 3}
        QIDCM = CSUM * 2.0 CO
        IF(IPT.GT.0) GO TO 11
        IF(IFLAG.EQ.1) RETURN
        WRITE(6,995) QIDCM,IFLAG,QERRCR
        FCMAT(1,15X,'RE:RESULTS:',E14.7,/,
          1.,15X,'I FLAG:',E14.7,/,
          2.,15X,'I3,2X,E14.7,/,',E14.7,/)
995    11

```

QLDR408E
QLDR409D
QLDR409S
QLDR4100
QLDR4105
QLDR4110
QLDR4115
QLDR4120
QLDR4125
QLDR4130
QLDR4135
QLDR4140
QLDR4145
QLDR4150
QLDR4155
QLDR4160
QLDR4165
QLDR4170
QLDR4175
QLDR4180
QLDR4185
QLDR4190
QLDR4195
QLDR4200
QLDR4205
QLDR4210
QLDR4215
QLDR4220
QLDR4225
QLDR4230
QLDR4235
QLDR4240
QLDR4245
QLDR4250
QLDR4255
QLDR4260
QLDR4265
QLDR4270
QLDR4275
QLDR4280
QLDR4285
QLDR4290
QLDR4295
QLDR4300
QLDR4305
QLDR4310
QLDR4315
QLDR4320

QLCR4565
QLCR4570
QLDR4575
QLDR4580
QLDR4585
QLDR4590
QLDR4595
QLDR4600
QLDR4605
QLDR4610
QLDR4615
QLDR4620
QLDR4625
QLDR4630
QLDR4635
QLDR4640
QLDR4645
QLDR4650
QLDR4655
QLDR4660
QLDR4665
QLDR4670
QLDR4675
QLDR4680
QLDR4685
QLDR4690
QLDR4695
QLDR4700
QLDR4705
QLDR4710
QLDR4715
QLDR4720
QLDR4725
QLDR4730
QLDR4735
QLDR4740
QLDR4745
QLDR4750
QLDR4755
QLDR4760
QLDR4765
QLDR4770
QLDR4775
QLDR4780
QLDR4785
QLDR4790
QLDR4795
QLDR4800

```

IF(IPT.LE.0) RETURN
CCTC 60
CIPRBP = CCPLX(C.CDO,0.0DO)
IF(IPT.LE.0) RETURN
WRITE(6,995) QIPRBP
995 FORMAT('0',10X,'QIPRBP = ',E14.7,' ',E14.7)
RETURN
END
COMPLEX FUNCTION CIPABP*16(DK,DR,CW,RHO,OFFSET,SIGMA,XSTN,IPT)
IMPLICIT REAL*8 (A-F,C,P,R-Y), COMPLEX*16 (Q,Z)
DIMENSION QINP(2)
IF(IPT.GT.0) WRITE(6,990) DK,DR,DW,RHO,OFFSET,SIGMA,XSTA,IPT
990 FCRMAT('C',10X,'QIPABP ENTERED WITH: ',/,
1. ,10X,'DK',11X,'DR',11X,'Dw',11X,'RHO',10X,'CFFSET',7),
2. ,SIGMA,18X,'XSTN',9X,'IPT',/,1,10X,7(E12.5,1,1),13)
XASTN = XSTN - OFFSET
IF(XASTN.LE.RHO-1.0DO) GOTO 20
ICF = IPT - 1
CCK = DCPLX(0.0DO,DK)
QCONST = CDEXP(DCPLX(0.0DO,SIGMA))
CALL QIPWXM(DK,DR,CW,RHC,XASTN,CINP,IOT)
QIPABP = (QCK*QINP(2) + QINP(1)) * QCONST
IF(IPT.LE.0) RETURN
GOTO 60
990 QIPABP = DCPLX(C.CDC,0.0DO)
IF(IPT.LE.0) RETURN
60 WRITE(6,995) QIPABP
995 FCRMAT('0',10X,'QIPABP = ',E14.7,' ',E14.7)
RETURN
END
SUBROUTINE CIPWXM(CK,DR,DW,RHO,XSTN,CINP,IPT)
N IS THE MAXIMUM DEGREE OF THE TERM U IN THE INTEGRALS
IMPLICIT REAL*8 (A-F,C,P,R-Y), COMPLEX*16 (Z,Q)
DIMENSION CINP(2)
IF(IPT.GT.0) WRITE(6,990) DK,DR,RHC,XSTN,IPT
990 FCRMAT('C',10X,'CIPWXM ENTERED WITH ARGUMENTS: ',/,
1. ,10X,'DK',16X,'DR',16X,'RHO',16X,'XSTN',14X,'IPT',/,
2. ,10X,4(E13.6,5X),12,5X,13)
IOT = IPT - 1
B = XSTN - RHO - 1.0C-10
A = -1.000
CALL CIPZXM(CK,DR,CW,RHO,XSTN,A,E,1,QANS,IOT)
CINP(1) = QANS
CALL CIPZXM(CK,DR,CW,RHC,XSTN,A,E,2,QANS,IOT)
CINP(2) = QANS
IF(IPT.LE.0) RETURN

```

CC

CLDR4805
CLDR4810
CLDR4815
CLDR4820
CLDR4825
CLDR4830
CLDR4835
CLDR4840
CLDR4845
CLDR4850
CLDR4855
CLDR4860
CLDR4865
CLDR4870
CLDR4875
CLDR4880
CLDR4885
CLDR4890
CLDR4895
CLDR4900
CLDR4905
CLDR4910
CLDR4915
CLDR4920
CLDR4925
CLDR4930
CLDR4935
CLDR4940
CLDR4945
CLDR4950
CLDR4955
CLDR4960
CLDR4965
CLDR4970
CLDR4975
CLDR4980
CLDR4985
CLDR4990
CLDR4995
CLDR5000
CLDR5005
CLDR5010
CLDR5015
CLDR5020
CLDR5025
CLDR5030
CLDR5035
CLDR5040

```

995 WRITE (6,995)
   FORMAT(' ',10X,'QIPWXM RESULTS:',/,
   ' ',11X,'J',6X,'QINP(I)',
   ' ',10X,'I',12X,'E12.6',/,E12.6)
996 CLDR4805
1001 WRITE (6,996) I, CIMP(I)
      RETURN
      END
SUBROUTINE QIPZXM (DK,DR,DW,RHO,XSTN,A,B,J,CANS,IPT)
  IMPLICIT REAL*8 (A - E,G,F,M,Q,P,R - Y), COMPLEX*16 (F,Q,Z)
  DIMENSION FV(5),LCFR(60),F1T(60),F2T(60),F3T(60),DAT(60),
  1 ARESTT(60),QUESTT(60),EPST(60),FSUM(60)
  1 F(X) = (X**J)*CDEXP(QEXP*(X))
  1 GAMMA = MMESJO((OMEGA * DSQRT((XSTN-X)*(XSTN-X) - YY)), IER)
  YY = RHO * RHO
  DM2 = (GAMMA + 1.0E0) * DW
  CMGMA = CSQRT (DK*CK*(1.0D0-DM2)/(CM2*DM2))
  CLAMDA = DK/CM2
  CM = DSQRT(CM2)
  C2EXP = CDEXP (DCMPLX (0.0D0, -DLAMDA*XSTN))
  QEXP = DCMPLX(0.0D0,CLAMDA)
  J1 = J - 1
  UACC = 9.0D-13
  IF (IPT.GT.0) WRITE(6,990)DK,DR,DW,RHO,XSTN,A,B,J,IPT
  990 1. ' ',15X,'DK',13X,'CR',13X,'DW',10X,'RHO',12X,'XSTN',11X,
  2. 'A',14X,'B',14X,'J',12X,'IPT',13X,
  3. ' ',15X,'7(E14.7',/,12,2X,I3),
      EFCURU = 4.0*U
  IFLAG = 1
  EFS = ACC
  CERROR = CCMPLX(C.0D0,0.0D0)
  LVL = 1
  LCRR(LVL) = 1
  CFSLM(LVL) = 0.0
  ALPHA = A
  CL1 = B - A
  AREA = 0.0
  AREST = 0.0
  FV(1) = F(ALPHA)
  FV(3) = F(ALPHA + 0.5*DA)
  FV(5) = F(ALPHA + CA)
  KCUNT = 3
  MWT = DA/6.0
  CEST = WT*(FV(1) + 4.0*FV(3) + FV(5))
  CDX = 0.5*DA
1

```

```

FV(2) = F(ALPHA + 0.5*DX)
FV(4) = F(ALPHA + 1.5*DX)
KCUNT = KCUNT + 2
WT = DX/6.0
QESTL = WT*(FV(1) + 4.0*FV(2) + FV(3))
QSUM = WT*(FV(3) + 4.0*FV(4) + FV(5))
QESTR = CESTL + QESTL
QESTL = WT*(CDABS(FV(1)) + CDABS(4.0*FV(2)) + CDABS(FV(3)))
ARESTL = WT*(CDABS(FV(3)) + CDABS(4.0*FV(4)) + CDABS(FV(5)))
AREA = AREA + ((ARESTL + ARESTR) - ARESTR)
CCIFF = QEST - QSUM
IF(CCABS(CCIFF).LE.EPS*CDABS(AREA))GO TO 2
IF(CCABS(CX).LE.EFOUR*CDABS(ALPHA))GO TO 5
IF(LVL.GE.60)GO TO 5
IF(KCUNT.GE.2000)GC TC 6
LVL = LVL + 1
LCRR(LVL) = 0
F1(LVL) = FV(3)
F2(LVL) = FV(4)
F3(LVL) = FV(5)
DATA = DX
DAT(LVL) = IX
ARESTL = ARESTL
ARESTR(LVL) = ARESTR
QESTL = CESTL
QESTR(LVL) = QESTR
CEFS = EPS/1.4
EFS(LVL) = EPS
FV(5) = FV(3)
FV(3) = FV(2)
GC TO 1
QERROR = QERROR + CDIFF/15.0
IF(LORR(LVL).EQ.0)GC TC 4
QSUM = QSUM + QESTL + QESTR
LVL = LVL - 1
IF(LVL.GT.1)GO TO 3
QANS = QSUM * QEXP / DM
IF(IPT.GT.0)GO TO 11
IF(IFLAG.EC.1)RETURN
WRITE(6,995) DK,DR,DM,RHO,XSIN,A,E,J,IPT
FCRMAT(1,995) QANS,IFLAG,IER,QERR,F,E14.7,/,
1. ,15X,IFLAG,2X,IER,5X,QERR,F,/,
2. ,15X,13,2X,13,5X,E14.7,/,E14.7,/,
RETURN
QSUM(LVL) = QSUM
LCRR(LVL) = 1
ALPHA = ALPHA + DA

```

```

CLDR5045
CLDR5050
CLDR5060
CLDR5070
CLDR5075
CLDR5080
CLDR5085
CLDR5090
CLDR5095
CLDR5100
CLDR5110
CLDR5115
CLDR5120
CLDR5130
CLDR5135
CLDR5140
CLDR5145
CLDR5150
CLDR5155
CLDR5160
CLDR5165
CLDR5170
CLDR5175
CLDR5180
CLDR5185
CLDR5190
CLDR5200
CLDR5210
CLDR5215
CLDR5220
CLDR5225
CLDR5230
CLDR5235
CLDR5240
CLDR5245
CLDR5250
CLDR5255
CLDR5260
CLDR5265
CLDR5270
CLDR5280

```



```

QIFRIP = -QANS
IF(IPT.LE.0) RETURN
GOTO 60
QIPRIP = DCMLPX(0.0D0,0.0D0)
2C IF(IPT.LE.0) RETURN
60 WRITE(6,995) QIPRIP
995 FORMAT('0',10X,'QIFRIP = ',E14.7,' ',E14.7)
RETURN
END
SUBROUTINE QIPZSM (DK,DR,DW,RHO,XSTN,A,B,J,CANS,QICF,IPT)
IMPLICIT REAL*8(A - E,G,H,M,O,P,R - Y), COMPLEX*16(F,Q,Z)
DIMENSION QICF(13)
DIMENSION FV(5),LORR(60),FIT(60),F2T(60),F3T(60),CAT(60),
1 F(X) = CLGNCR(X,CR,J,QICF)*CDEXP(CEXP*(X))
1 GAMMA = 1.4E0
1 YV = RHC * RHO
DM2 = (GAMMA + 1.0D0) * DW
OMEGA = DSQRT(DK*DK*(1.0D0-DM2)/(DM2*DM2))
CLAMDA = DK/DM2
DM = DSQRT(DM2)
QZEXP = CDEXP(DCMLPX(0.0D0,-CLAMDA*XSTN))
QEXP = DCMLPX(0.0D0,CLAMDA)
UACC = 9.0D-13
IF(IPT.GT.0) WRITE(6,990)DK,DR,DH,RHO,XSTN,A,B,J,IPT
990 FCFORMAT('0',15X,'QIPZSM ENTERED WITH ARGUMENTS:',/,
1.,15X,'DK',13X,'DR',13X,'DW',10X,'RHC',12X,'XSTN',11X,
2.,14X,'B',14X,'J',12X,'IPT',/,
3.,15X,'7(E14.7,.,.),12,2X,I3)
EFOURU = 4.*U
IFLAG = 1
EPS = ACC
CERROR = CCMLPX(0.0D0,0.0D0)
LV1 = 1
LCRR(LV1) = 1
QFSUM(LV1) = 0.0
ALPHA = A
DA = B - A
AREA = 0.0
AFEST = 0.0
FV(1) = F(ALPHA)
FV(3) = F(ALPHA) + C.5*DA
FV(5) = F(ALPHA) + DA
KCOUNT = 3
WT = DA/6.0
CEST = WT*(FV(1) + 4.0*FV(3) + FV(5))

```

QLDR55250
 QLDR55255
 QLDR55260
 QLDR55265
 QLDR55270
 QLDR55275
 QLDR55280
 QLDR55285
 QLDR55290
 QLDR55295
 QLDR55300
 QLDR55305
 QLDR55310
 QLDR55315
 QLDR55320
 QLDR55325
 QLDR55330
 QLDR55335
 QLDR55340
 QLDR55345
 QLDR55350
 QLDR55355
 QLDR55360
 QLDR55365
 QLDR55370
 QLDR55375
 QLDR55380
 QLDR55385
 QLDR55390
 QLDR55395
 QLDR55400
 QLDR55405
 QLDR55410
 QLDR55415
 QLDR55420
 QLDR55425
 QLDR55430
 QLDR55435
 QLDR55440
 QLDR55445
 QLDR55450
 QLDR55455
 QLDR55460
 QLDR55465
 QLDR55470
 QLDR55475
 QLDR55480
 QLDR55485
 QLDR55490
 QLDR55495
 QLDR55500
 QLDR55505
 QLDR55510
 QLDR55515
 QLDR55520
 QLDR55525
 QLDR55530
 QLDR55535
 QLDR55540
 QLDR55545
 QLDR55550
 QLDR55555
 QLDR55560
 QLDR55565
 QLDR55570
 QLDR55575
 QLDR55580
 QLDR55585
 QLDR55590
 QLDR55595
 QLDR55600
 QLDR55605
 QLDR55610
 QLDR55615
 QLDR55620
 QLDR55625
 QLDR55630
 QLDR55635
 QLDR55640
 QLDR55645
 QLDR55650
 QLDR55655
 QLDR55660
 QLDR55665
 QLDR55670
 QLDR55675
 QLDR55680
 QLDR55685
 QLDR55690
 QLDR55695
 QLDR55700
 QLDR55705
 QLDR55710
 QLDR55715
 QLDR55720
 QLDR55725
 QLDR55730
 QLDR55735
 QLDR55740
 QLDR55745
 QLDR55750
 QLDR55755
 QLDR55760

Q1CR5765
Q1CR5770
Q1CR5775
Q1CR5780
Q1CR5785
Q1CR5790
Q1CR5795
Q1CR5800
Q1CR5805
Q1CR5810
Q1CR5815
Q1CR5820
Q1CR5825
Q1CR5830
Q1CR5835
Q1CR5840
Q1CR5845
Q1CR5850
Q1CR5855
Q1CR5860
Q1CR5865
Q1CR5870
Q1CR5875
Q1CR5880
Q1CR5885
Q1CR5890
Q1CR5895
Q1CR5900
Q1CR5905
Q1CR5910
Q1CR5915
Q1CR5920
Q1CR5925
Q1CR5930
Q1CR5935
Q1CR5940
Q1CR5945
Q1CR5950
Q1CR5955
Q1CR5960
Q1CR5965
Q1CR5970
Q1CR5975
Q1CR5980
Q1CR5985
Q1CR5990
Q1CR5995
Q1CR6000

```

1  DX = 0.5*CA
   FV(2) = F(ALPHA + 0.5*DX)
   FV(4) = F(ALPHA + 1.5*DX)
   KOUNT = KOUNT + 2
   WT = DX/6.0
   QESTL = WT*(FV(1) + 4.0*FV(2) + FV(3))
   QSUM = CESTL + QESTR
   ARESTL = WT*(CDABS(FV(1)) + CDABS(4.0*FV(2)) + CDABS(FV(3)))
   AREA = AREA + ((ARESTL + ARESTR) - AREST)
   QDIFF = QEST - QSUM
   IF(CDABS(QDIFF).LE.EPS*CDABS(AREA))GO TO 2
   IF(DABS(CX).LE.EFCURU*DABS(ALPHA))GO TO 5
   IF(LVL.GE.6)GO TO 5
   IF(KOUNT.GE.2000)GC TC 6
   LVL = LVL + 1
   LCRR(LVL) = 0
   F1(LVL) = FV(3)
   F2(LVL) = FV(4)
   F3(LVL) = FV(5)
   DA = DX
   LAT(LVL) = DX
   ARESTL = ARESTL + ARESTR
   ARESTT(LVL) = ARESTR
   CESTT(LVL) = CESTL
   QESTT(LVL) = QESTR
   EPS = EPS/1.4
   EPST(LVL) = EPS
   FV(5) = FV(2)
   FV(3) = FV(2)
   GC TO 1
   GEFFGR(LVL) = QERROR + QDIFF/15.0
   IF(LORR(LVL).EQ.0)GO TO 4
   QSUM = QPSUM(LVL) + QSUM
   LVL = LVL - 1
   IF(LVL.GT.1)GO TO 3
   QANS = QSUM * Q2EXP / DM
   IF(IPT.GT.0)GC TC 11
   IF(IFLAG.EQ.1)RETURN
   WRITE(6,995)DK,DR,DM,RHO,XSTN,A,E,J,IPT
   IF(ITE(6,995)QANS,IFLAG,IER,ERRCR,/,
995 1. ,15X,13,2X,13,5X,E14.7,/,
2 1. ,15X,13,2X,13,5X,E14.7,/,
   RETURN
   QPSUM(LVL) = QSUM
   LCRR(LVL) = 1

```


QLDR6245
 QLDR6250
 QLDR6255
 QLDR6260
 QLDR6265
 QLDR6270
 QLDR6275
 QLDR6280
 QLDR6285
 QLDR6290
 QLDR6295
 QLDR6300
 QLDR6305
 QLDR6310
 QLDR6315
 QLDR6320
 QLDR6325
 QLDR6330
 QLDR6335
 QLDR6340
 QLDR6345
 QLDR6350
 QLDR6355
 QLDR6360
 QLDR6365
 QLDR6370
 QLDR6375
 QLDR6380
 QLDR6385
 QLDR6390
 QLDR6395
 QLDR6400
 QLDR6405
 QLDR6410
 QLDR6415
 QLDR6420
 QLDR6425
 QLDR6430
 QLDR6435
 QLDR6440
 QLDR6445
 QLDR6450
 QLDR6455
 QLDR6460
 QLDR6465
 QLDR6470
 QLDR6475
 QLDR6480

```

55C FCRMAT('0',10X,'QICRBP ENTERED WITH:','/,
1  ',10X,'OK',11X,'CR',11X,'DW',11X,'RHC',1C),XSTN',9X,'IPT',
2  ',10X,5(E12.5,'),13)
  IF(XSTN.LE.RHO-1.000) GOTO 20
  ICT = IPT - 1
  QCK = DCMPLX(0.000,DK)
  CALL QIDWXM(DK,DR,DW,FHO,XSTN,QINP,IOT)
  QIDRBP = -QCK*QINP(2) - QINP(1)
  IF(IPT.LE.0) RETURN
  GCIC 60
20 QIDRBP = DCMPLX(0.000,0.000)
  IF(IPT.LE.0) RETURN
60 WRITE(6,995) QIDREP
995 FCRMAT('0',10X,'QICRBP = ',E14.7,' ',E14.7)
  RETURN
ENCL
C COMPLEX FUNCTION QIDABP*16(DK,DR,DW,RHO,CFFSET,SIGMA,XSTN,IPT)
  IMPLICIT REAL*8 (A-F,O,P,R-Y), COMPLEX*16 (Q,Z)
  DIMENSION QINP(2)
  IF(IPT.GT.0) WRITE(6,990) DK,CR,CW,RHO,OFFSET,SIGMA,XSTN,IPT
55C FCRMAT('C',10X,'QICABP ENTERED WITH:','/,
1  ',10X,'OK',11X,'CR',11X,'DW',11X,'RHO',10X,'OFFSET',7X,
2  ',SIGMA',8X,'XSTN',5X,'IPT',/,',10X,7(E12.5,'),13)
  XSTN = XSTN - OFFSET
  IF(XSTN.LE.RHO-1.000) GOTO 20
  ICT = IPT - 1
  QCK = DCMPLX(0.000,DK)
  CCONST = CDEXP( DCMPLX(0.000,SIGMA))
  CALL QIDWXM(DK,DR,DW,RHO,XASTN,QINP,IOT)
  QIDABP = (QCK*QINP(2) + QINP(1)) * QCONST
  IF(IPT.LE.0) RETURN
  GC TO 60
20 QIDABP = DCMPLX(0.000,0.000)
  IF(IPT.LE.0) RETURN
60 WRITE(6,995) QIDABP
995 FCRMAT('0',10X,'CIDABP = ',E14.7,' ',E14.7)
  RETURN
END
SUBROUTINE QIDWXM(CK,DR,DW,RHO,XSTN,QINP,IPT)
C
C
  A IS THE MAXIMUM DEGREE OF THE TERM U IN THE INTEGRALS
  IMPLICIT REAL*8 (A-F,C,P,R-Y), COMPLEX*16 (Z,C)
  DIMENSION QINP(2)
  IF(IPT.GT.0) WRITE(6,990) DK,CR,RHC,XSTN,IPT
55C FCRMAT('0',10X,'QIDWXM ENTERED WITH ARGUMENTS:',
1  ',10X,'OK',16X,'DR',16X,'RHO',15X,'XSTN',14X,'IPT',/,
2  ',10X,4(E13.6,5X),12,5X,13)
  
```


Q1DR67220
 Q1DR67230
 Q1DR67240
 Q1DR67250
 Q1DR67260
 Q1DR67270
 Q1DR67280
 Q1DR67290
 Q1DR67300
 Q1DR67310
 Q1DR67320
 Q1DR67330
 Q1DR67340
 Q1DR67350
 Q1DR67360
 Q1DR67370
 Q1DR67380
 Q1DR67390
 Q1DR67400
 Q1DR67410
 Q1DR67420
 Q1DR67430
 Q1DR67440
 Q1DR67450
 Q1DR67460
 Q1DR67470
 Q1DR67480
 Q1DR67490
 Q1DR67500
 Q1DR67510
 Q1DR67520
 Q1DR67530
 Q1DR67540
 Q1DR67550
 Q1DR67560

```

    AFEA = 0.0
    AREST = 0.0
    FV(1) = F(ALPHA)
    FV(3) = F(ALPHA + 0.5*CA)
    FV(5) = F(ALPHA + CA)
    KCUNT = 3
    WT = DA/6.0
    QEST = WT*(FV(1) + 4.0*FV(3) + FV(5))
    DX = 0.5*DA
    FV(2) = F(ALPHA + 0.5*DX)
    FV(4) = F(ALPHA + 1.5*DX)
    KCUNT = KCUNT + 2
    WT = DX/6.0
    I STL = WT*(FV(1) + 4.0*FV(2) + FV(3))
    CESTR = WT*(FV(3) + 4.0*FV(4) + FV(5))
    CSLM = QESTL + QESTR
    ARESTL = WT*(CDABS(FV(1)) + CDABS(4.0*FV(2)) + CDABS(FV(3)))
    ARESTR = WT*(CDABS(FV(3)) + CDABS(4.0*FV(4)) + CDABS(FV(5)))
    AFEA = AFEA + ((ARESTL + ARESTR) - AREST)
    QDIFF = QEST - QSUM
    IF(CDABS(QDIFF).LE.EPCURU*DABS(ALPHA))GO TO 2
    IF(LVL.GE.6)GO TO 5
    IF(KCUNT.GE.400)GC TC 6
    LVL = LVL + 1
    LCRR(LVL) = 0
    FIT(LVL) = FV(3)
    FET(LVL) = FV(4)
    FET(LVL) = FV(5)
    LA = DX
    LA(T(LVL)) = CX
    ARESTL = ARESTL
    ARESTT(LVL) = ARESTR
    CESTL = CESTL
    QESTT(LVL) = QESTR
    QESTS = EPS/1.4
    EPS = EPS
    FV(5) = FV(3)
    FV(3) = FV(2)
    GC TO 1
    QERRCR = QERROR + QDIFF/15.0
    IF(LORR(LVL).EC.0)GC TO 4
    QSUM = CPSUM(LVL) + QSUM
    LVL = LVL - 1
    IF(LVL.GT.1)GO TO 3
    QANS = ((B*JI)*CDEXP(QEXP*B) - (CSLM)*Q2EXP/DW
    IF(IPT.GT.0)GO TC 11
    IF(FLAG.EQ.1) RETURN
  
```

1

2 3

109

Q1 DR7205
Q1 CR7210
Q1 CR7215
Q1 DR7225
Q1 CR7225
Q1 CR7230
Q1 CR7235
Q1 DR7240
Q1 DR7245
Q1 CR7250
Q1 DR7255
Q1 CR7260
Q1 CR7265
Q1 CR7270
Q1 CR7275
Q1 CR7280
Q1 CR7290
Q1 CR7295
Q1 CR7300
Q1 CR7305
Q1 CR7310
Q1 CR7315
Q1 CR7320
Q1 CR7325
Q1 CR7330
Q1 CR7335
Q1 CR7340
Q1 CR7345
Q1 CR7350
Q1 CR7355
Q1 CR7360
Q1 CR7365
Q1 CR7370
Q1 CR7375
Q1 CR7380
Q1 CR7390
Q1 CR7395
Q1 CR7400
Q1 CR7405
Q1 CR7410
Q1 CR7415
Q1 CR7420
Q1 CR7425
Q1 CR7430
Q1 CR7435
Q1 CR7440

```

ALPHA = A
DA = B - A
AREA = 0.0
AREST = 0.0
FV(1) = F(ALPHA) + 0.5*DA
FV(3) = F(ALPHA + DA)
FV(5) = F(ALPHA + DA)
KKCNT = 3
WT = DA/6.0
CEST = WT*(FV(1) + 4.0*FV(3) + FV(5))
DX = 0.5*DA
FV(2) = F(ALPHA + C.5*DX)
FV(4) = F(ALPHA + 1.5*DX)
KVCNT = KOUNT + 2
WTLSTL = DX/6.0
QESTL = WT*(FV(1) + 4.0*FV(2) + FV(3))
CESTR = WT*(FV(3) + 4.0*FV(4) + FV(5))
QSUM = QESTL + CESTR
QAFCSTL = WT*(CDABS(FV(1))) + CDABS(4.0*FV(2)) + CDABS(FV(3)))
AFCRESTR = WT*(CDABS(FV(3))) + CDABS(4.0*FV(4)) + CDABS(FV(5)))
AREA = AREA + ((ARESTL + ARESTR) - AREST)
CCIFF = QEST - QSUM
IF(CDABS(QDIFF).LE.EPS*DABS(AREA))GO TO 2
IF(DABS(OX).LE.EFCURU*DABS(ALPHA))GO TO 5
IF(LVL.GE.6)GO TIC 5
IF(KCUNT.GE.4000)GC TO 6
LVL = LVL + 1
LCRR(LVL) = 0
FFIT(LVL) = FV(3)
FFET(LVL) = FV(4)
FFET(LVL) = FV(5)
DATA(LVL) = DX
CAFCSTL = ARESTL
QAFCSTL(LVL) = ARESTR
CESTT(LVL) = QESTL
CESTRT(LVL) = CESTR
EPS = EPS/1.4
EEFV(5) = FV(3)
EEFV(3) = FV(2)
EEFV(TIC 1) = FV(TIC 1)
ERRCR = CEROR + CCIFF/15.0
IF(LCRR(LVL).EQ.0)GO TO 4
CSUM = QPSUM(LVL) + CSUM
LVL = LVL - 1
IF(LVL.GT.1)GO TO 3
IGANS = (CLGNDR(B,IR,J,QICF)*CDEXF(CEXP*B) - QSUM)*QZEXP/DM

```

1

23

7. Gorelov Spanning Function.

The subprogram for function used by Gorelov is shown below.

```
REAL FUNCTION PLGNDR*8(X,DR,N)
IMPLICIT REAL*8(A-H,O-Z)
IF(N.EQ.0) GOTO 100
ETA = CARCCS(-X)
ETASTR = DARCOS(1.000 - DR)
FN = CFLCAT(N)
PLGNDR = DCCS(FN*ETA)-DCQS(FN*ETASTR)
RETURN
100 PLGNDR = 1.000
RETURN
END
```


8. Linear Expansion Program

The program based on the linear expansion for small k is shown below.

LIN00490
 LIN00500
 LIN00510
 LIN00520
 LIN00530
 LIN00540
 LIN00550
 LIN00560
 LIN00570
 LIN00580
 LIN00590
 LIN00600
 LIN00610
 LIN00620
 LIN00630
 LIN00640
 LIN00650
 LIN00660
 LIN00670
 LIN00680
 LIN00690
 LIN00700
 LIN00710
 LIN00720
 LIN00730
 LIN00740
 LIN00750
 LIN00760
 LIN00770
 LIN00780
 LIN00790
 LIN00800
 LIN00810
 LIN00820
 LIN00830
 LIN00840
 LIN00850
 LIN00860
 LIN00870
 LIN00880
 LIN00890
 LIN00900
 LIN00910
 LIN00920
 LIN00930
 LIN00940
 LIN00950
 LIN00960

```

END SUBROUTINE ABSA (A,OFFSET,X,DR,DW,CCOMPLEX*16(Q,Z))
  IMPLICIT REAL*8 (A-H,C,P,R-Y), COMPLEX*16(Z,Q)
  DIMENSION X(13)
  XINT = (2.0D0-DR)/CFLOAT(N+1)
  XL = CR - 1.0D0 + 1.0D-8
  DO 10 I = 1,N
    XI(I) = XL + (XINT*CFLOAT(I))
    IF(X(I).EQ.0.0D0) X(I) = .1D-14
  CONTINUE
  RETURN
10
END
SUBROUTINE Q$ONE(DK,DR,DW,RHO,OFFSET,SIGMA,N,NF,IPT,X)
  IMPLICIT REAL*8(A-H,C,P,R-Y), COMPLEX*16(Z,Q)
  DIMENSION QICOF(26), QIINT(26,26), QIRBPF(13), QIABP(13)
  DIMENSION QIRBCF(13), QIABCF(13)
  IF(IPT.GT.0) WRITE(6,990)DK,DR,RHO,OFFSET,SIGMA,N,NF,IPT
  ICT = IPT - 1
  FCFORMAT(0) SUBROUTINE Q$ONE ENTERED WITH: ',',
1. 5X,DK,13X,CR,13X,RHO,12X,OFFSET,9X,SIGMA,10X,
2. N,3X,NF,2X,IPT,1.
3. 5X,5(E12.5,3X),12,2X,13)
  CALL QIZINT(DK,DR,DW,RHO,OFFSET,SIGMA,N,X,ICT,QICCF,QIINT)
  CALL QICCOF(QICOF,QIINT,N,IPT,QIRBCF,QIRBCF)
  CALL QIPOT(DK,DR,DW,RHO,OFFSET,SIGMA,N,X,QIAECF,QIRBCF,ICT)
  IQT = 0
  CALL QICCOF(DK,DR,DW,RHO,OFFSET,SIGMA,N,QIAECF,QIRBCF,IQT)
  RETURN
END
CC
SUBROUTINE QIZINT(DK,DR,DW,RHO,OFFSET,SIGMA,N,X,
  IPT,QICOF,QIINT)
  IMPLICIT REAL*8(A-H,C,P,R-Y), COMPLEX*16(Z,Q)
  DIMENSION QICOF(26), QIINT(26,26), QINTRP(13), QINTAP(13)
  DIMENSION X(13)
  FCFORMAT(0) QIZINT ENTERED WITH: ',',
1. 10X,DK,13X,DR,13X,RHO,12X,OFFSET,9X,SIGMA,10X,
2. N,2X,IPT,1.
3. 10X,5(E12.5,3X),12,2X,13)
  CCONST = 5(DEXP(0.0D0,SIGMA))
  IF(IPT.GT.0) WRITE(6,990) DK,DR,RHO,OFFSET,SIGMA,N,IPT
  ICT = IPT - 1
  GAMMA = 1.4D0
  CLAMDA = DK/((GAMMA+1.0D0)*DW)
  DO 90 I = 1,N
990
  
```

INCG57C
 LIN0058C
 LIN0059C
 LIN0100C
 LIN0101C
 LIN0102C
 LIN0103C
 LIN0104C
 LIN0105C
 LIN0106C
 LIN0107C
 LIN0108C
 LIN0109C
 LIN0110C
 LIN0111C
 LIN0112C
 LIN0113C
 LIN0114C
 LIN0115C
 LIN0116C
 LIN0117C
 LIN0118C
 LIN0119C
 LIN0120C
 LIN0121C
 LIN0122C
 LIN0123C
 LIN0124C
 LIN0125C
 LIN0126C
 LIN0127C
 LIN0128C
 LIN0129C
 LIN0130C
 LIN0131C
 LIN0132C
 LIN0133C
 LIN0134C
 LIN0135C
 LIN0136C
 LIN0137C
 LIN0138C
 LIN0139C
 LIN0140C
 LIN0141C
 LIN0142C
 LIN0143C
 LIN0144C

```

IN = I + N
XSTN = X(I) - DR
XCEP = CDEXP(DCMPLX(C.000, DLAMDA*XSTN))
CALL QICARP(DK, DR, DW, RHO, OFFSET, SIGMA, XSTN, CINTAP, N, IOT)
CALL QICRIP(DK, DR, DW, RHC, OFFSET, XSTN, QINTRP, N, ICT)
DO 20 J = 1, N
  JN = N + J
  J1 = J-1
  J1INT(I, J1) = DCMPLX(1.000, CLAMDA*XSTN)
  TEMP = CR-XSTN-1.000
  Q1INT(I, 2) = DCMPLX(TEMP, CLAMCA*XSTN*TEMP)
  Q1INT(IN, J) = QINTRP(J)
  Q1INT(I, JN) = QINTAP(J)
  XASTN = XSTN - OFFSET
  Q1INT(IN, 3) = DCMPLX(1.000, DLAMDA*XASTN)
  TEMP = DR-XASTN-1.000
  Q1INT(IN, 4) = DCMPLX(TEMP, DLAMDA*XASTN*TEMP)
  CCNTINUE
20  Q1CCF(IN) = -QICRBP(DK, DR, DW, RHO, XSTN, IOT)
  Q1CCF(I) = -QICRBP(DK, DR, DW, RHC, CFFSET, SIGMA, XSTN, IOT)
  CCNTINUE
90  FCTURN
  ENCL
  COMPLEX FUNCTION QICRBP*16(DK, DF, DW, RHO, XSTN, IPT)
  IMPLICIT REAL*8 (A-H, O, P, R-Y), COMPLEX*16 (Q, Z)
  DIMENSION QINP(2)
  IF (IPT.GT.0) WRITE(6, 990) CK, CR, CW, RHO, XSTN, IPT
  IF (IPT.LE.0) CK, CR, CW, RHO, XSTN, IPT
  FCFORMAT(0, 10X, 'QICRBP ENTERED WITH: ', /,
    1, 10X, 'DK, 11X, 'CR, 11X, 'DW, 11X, 'RHO, 10X, 'XSTN, 9X, 'IPT,
    2, /, 'XSTN, 1E, RHO-1.000) GOTC 20
  IF (XSTN.GT.2.000) GOTC 20
  IF (IPT - 1) GOTC 20
  ICT = DCMPLX(0.000, CK)
  GAMMA = 1.400
  CLAMDA = DK/((GAMMA+1.000)*DW)
  QICRBP = DCMPLX(1.000, (DK+DLAMDA)*(XSTN-DR))
  IF (IPT.LE.0) RETURN
  GOTC 30
  QICRBP = CCMPLX(0.000, 0.000)
  IF (IPT.LE.0) RETURN
  WRITE(6, 995) QICRBP
995 FORMAT('C, 10X, 'QICRBP = ', E14.7, ', ', E14.7)
  RETURN
  ENCL
  COMPLEX FUNCTION QICABP*16(CK, DR, CW, RHO, OFFSET, SIGMA, XSTN, IPT)
  IMPLICIT REAL*8 (A-H, O, P, R-Y), COMPLEX*16 (Q, Z)

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DIMENSION QINP(2)
IF(IPT.GT.0) WRITE(6,990) DK,DR,DW,RHO,OFFSET,SIGMA,XSTN,IPT
FORMAT(10,'10X',QICABP,ENTERED WITH:','/,RHO,'10X',OFFSET,'7X',
'10X',11X,'DR',11X,'DW',11X,'10X',7(E12.5,' '),13)
'SIGMA',8X,'XSTN',9X,IPT,'/',10X,7(E12.5,' '),13)
XSTN=XSTN-OFFSET
IF(XASTN.LE.RHO-1.0D0) GOTO 20
IF(XASTN.GT.2.0D0) GOTO 20
IPT=IPT-1
CLK=DCMLPX(0.0D0,DK)
CCONST=CDEXP(CMLPX(0.0D0,SIGMA))
GAMMA=1.4D0
CLAMDA=CLK/((GAMMA+1.0D0)*DW)
QICABP=DCMLPX(1.0D0,(DK+CLAMDA))*((XASTN-CR))*QCONST
IF(IPT.LE.0) RETURN
GO TO 30
20 QICABP=DCMLPX(C.0D0,0.0D0)
IF(IPT.LE.0) RETURN
30 WRITE(6,995) QICABP
995 FORMAT(10,'10X',QICABP=' ',E14.7,' ',E14.7)
END
SUBROUTINE QICAIIP(CLK,DR,DW,RHO,CFST,SGMA,XSTN,QCAIP,N,IPT)
IMPLICIT REAL*8 (A-H,C,P,R-Y), COMPLEX*16 (C,Z)
DIMENSION QINP(13),QCAIP(13)
IF(IPT.GT.0) WRITE(6,990) CLK,CR,CW,RHO,CFST,SGMA,XSTN,IPT
IF(IPT.GT.0) WRITE(6,990) CLK,CR,CW,RHO,CFST,SGMA,XSTN,IPT
FCRMAT(0,'10X',11X,'DR',11X,'DW',11X,'10X',7(E12.5,' '),13)
'10X',11X,'10X',9X,IPT,'/',10X,7(E12.5,' '),13)
XSTN=XSTN-CFST
IF(XASTN.LE.DR+RHC-1.0D0) GO TO 20
IF(XASTN.GT.2.0D0) GOTO 20
IPT=IPT-1
CCONST=CDEXP(CMLPX(0.0D0,SGMA))
GAMMA=1.4D0
CLAMDA=CLK/((GAMMA+1.0D0)*DW)
QCAIP(1)=DCMLPX(1.0D0,CLAMDA*(XASTN-RHC))
TEMP=CR+RHO-XASTN-1.0D0
QCAIP(2)=DCMLPX(TEMP,CLAMDA*(XASTN-RHO)*TEMP)
IF(IPT.LE.0) RETURN
GOTO 30
20 ZERO=DCMLPX(0.0D0,0.0D0)
ZC 25 I=1,N
QCAIP(I)=ZERO
CONTINUE
IF(IPT.LE.0) RETURN
30 WRITE(6,995)
995 FORMAT(10,'10X',QICAIIP RESULTS: J QCAIP(J)')

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92 WRITE(6,92) I, I, QICOF(I)
   FORMAT('0',10X,'QICDEF EQUATION SYSTEM, ROW ',I2,'/',
   ' ',10X,'QICOF(',I2,',',E14.7,',',E14.7)
   CC 2 J = 1,N
   J2=J+N
93 WRITE(6,91) I,J,QIINT(I,J),I,J2,QIINT(I,J2)
   FORMAT(' ',15X,2('QINT(',I2,',',I2,',') = ',E14.7,',',E14.7,10X))
94 CC CONTINUE
   I 1 IA = 26
   I 2 IJOB=0
   CALL LECTIC(QIINT,A2,IA,QICOF,M,IE,IJOB,ZWA,IER)
   IF(IER.EQ.0) GOTC 30
   IF(IER.EQ.129) GC TO 10
   WRITE(6,93)
95 FCFORMAT('0',10X,'CICDEF - ITERATIVE IMPROVEMENT FAILED, MATRIX TOO
   ILL-CONDITIONED. USE RESULTS WITH CAUTION.')
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WRITE (6,95) I, XSTN
WRITE (6,91) COR, QCRP, CDAP, QCRIPQ, CCAIPP
WRITE (6,92) CDA, QCRP, CDAO, QDRIPP, CCAIPO
QCRP(I) = QCR
QCAA(I) = QCA
FCRMAI(1) = 1.0 X-STATIC NUMBER, I2, XSTN = F6.4, /,
1 BL TOTAL D(POT)/DX, 14X, REF EL D(POT)/DX, 8X,
2 ADJ BL D(POT)/DX, 8X, REF BL INT D(POT)/DX, 5X,
3 ADJ BL INT D(POT)/DX,
1C CONTINUE
954 WRITE(6,954)
1C FCRMAT(1,1,10X, 'SUMMARY LISTING', /,
1C, 10X, XSTN, 7X, 'SINGLE BLADE TOTAL POTENTIAL', 6X,
1 REF ELADE POTENTIAL, 15X, ADJ ELADE POTENTIAL,
1 DC 20 I = 1, N
XSTN = X(I)
WRITE(6,94) XSTN, CPHI(I), QRR(I), CAA(I)
94 FCRMAT(1,1,10X, F6.4, 3(5X, E14.7, , E14.7))
20 CONTINUE
956 WRITE(6,996)
1C FCRMAT(1,0,10X, XSTN, 7X, 'SINGLE ELADE TOTAL C(POT)/CX', 6X,
1 REF ELADE D(POT)/DX, 15X, ADJ ELADE D(POT)/DX,
1 CC 50 I = 1, N
XSTN = X(I)
WRITE(6,94) XSTN, CPHI(I), QCRP(I), CDA(I)
5C CONTINUE
END
SUBROUTINE QDCOF(CK, DR, DW, RHO, CFFSET, SIGMA, N, QIABCF, QIRBCF, IPT)
IMPLICIT REAL * 8 (A-F, Q, P, R-Y), COMPLEX * 16 (Q, Z)
DIMENSION QIABCF(13), QIRBCF(13)
IF(IPT, GT, 0) WRITE(6,990) DK, DR, CK, RHO, OFFSET, SIGMA, N, IPT
FCRMAI(1,1,10X, QDCOF - CALCULATION OF COMPLEX DIMENSIONLESS AERC
1 DYNAMIC CCEFFICIENTS, /,
3 0, 10X, DK, 13X, DR, 13X, DW, 13X, RHO, 12X, OFFSET, 5X, SIGMA,
4 10X, 10X, 3X, IPT, /,
5 (E12.5, 3X), 12, 2X, 13)
ICT = IPT - 3
QCLL = QCLL(CK, DR, DW, RHO, OFFSET, SIGMA, N, QIABCF, QIRBCF, IPT)
QCCM = QDCM(CK, DR, DW, RHO, OFFSET, SIGMA, N, QIABCF, QIRBCF, IPT)
GAMMA = 1.4CQ
TAU = (2.0DQ*DSQRT((GAMMA + 1.0CQ)*DW))/DR
WRITE(6,90) DK, TAU, DW, N, SIGMA, CCL, QDCM
90 FCRMAT(1,0,5X, DK = F6.3, , TAU = F7.4, , DW = F6.3, , N =
1 I2, , SIGMA = F6.3, , CL = F9.4, , CM = ,
2 F9.4, , F9.4)
END

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LI N04800

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ARESTR = WT*(CDABS(FV(3)) + CCABS(4.0*FV(4)) + COABS(FV(5)))
AREA = ARESTR - QSUM
QDIFF = ARESTR - ARESTR
IF (CCABS(QDIFF).LE.EPS*DABS(AREA)) GO TO 2
IF (DABS(DX).LE.EFCURU*DABS(ALPHA)) GO TO 5
IF (LVL.GE.60) GO TO 5
IF (KCUNT.GE.4000) GO TO 6
LVL = LVL + 1
LCRR(LVL) = 0
FIT(LVL) = FV(3)
F2T(LVL) = FV(4)
F3T(LVL) = FV(5)
CA = CX
DAT(LVL) = DX
ARESTR = ARESTR
ARESTR(LVL) = ARESTR
QEST = QESTL = QESTR
QESTT(LVL) = QESTR
EPS = EPS/1.4
EPST(LVL) = EPS
FV(1) = FV(3)
FV(2) = FV(4)
FV(3) = FV(5)
GC TO 1
CERROR = QERROR + QDIFF/15.0
IF (LCRR(LVL).EQ.0) GO TO 4
QSUM = QSUM(LVL) + CSUM
LVL = LVL - 1
IF (LVL.GT.1) GO TO 3
QIDCL = CSUM*2.0E0
IF (IPT.GT.0) GO TO 11
IF (IFLAG.EQ.1) RETURN
WRITE(6,995) QIDCL,IFLAG,CERROR
FCRMAT(1,15X,RESULTS:QIDCL = ,E14.7,,E14.7,,
1. ,15X,IFLAG,2X,QERROR,/,
2. ,15X,13,2X,E14.7,,E14.7)
RETURN
QSUM(LVL) = QSUM
LCRR(LVL) = 1
ALPHA = ALPHA + DA
DA = DAT(LVL)
FV(1) = FIT(LVL)
FV(3) = F2T(LVL)
FV(5) = F3T(LVL)
ARESTR = ARESTR(LVL)
EPS = EPST(LVL)
GC TO 1

```

2 3

11 995

4


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60 WRITE (6,995) C1FFBP
995 FCRMAT(0,10X,'Q1PRBP = ',E14.7,',' ,E14.7)
RETURN
ENC
COMPLEX FUNCTION Q1PAEP*16(DK,DR,DW,RHO,CFFSET,SIGMA,XSTN,IPT)
IMPLICIT REAL*8 (A-H,O,P,R-Y), COMPLEX*16 (Q,Z)
DIMENSION QINP(2)
IF(IPT,GT,0) WRITE(6,990) DK,DR,DW,RHO,OFFSET,SIGMA,XSTN,IPT
FORMAT(0,10X,'Q1PABP ENTERED WITH:',/
1 10X,'DK:',11X,'DR:',11X,'DW:',11X,'RHO:',10X,'CFFSET',7X,
2 10X,'SIGMA:',8X,'XSTN:',9X,'IPT',/,',10X,7(E12.5,',' ,13)
XSTN=XSTN-OFFSET
IF(XSTN.LE.RHO-1.0D0) GOTQ 20
ICT=IPT-1
CCNST=CDEXP(DCMPLX(0.0D0,SIGMA))
GAMMA=1.4D0
DM2=(GAMMA+1.0D0)*DW
DLAMDA=DK/DM2
PR=1.0D0+XSTN-RHO
FIMAG=(DK+DLAMDA)*((XSTN-RHO)*(XSTN-RHO)-1.0D0)
Q1PAEP=DCMPLX(PR,PI*FIMAG-DLAMDA*XSTN*PR)*CCNST/DSCRT(DM2)
IF(IPT,LE,0) RETURN
GOTO 60
20 Q1PAEP=CCMPLX(0.0D0,0.0D0)
IF(IPT,LE,0) RETURN
60 WRITE (6,995) Q1FABP
995 FCRMAT(0,10X,'Q1PABP = ',E14.7,',' ,E14.7)
RETURN
ENC
COMPLEX FUNCTION C1OPHI*16 (DK,DR,DW,RHO,CFFSET,SIGMA,
N,Q1ABCF,Q1RBCF,XSTN,IPT)
IMPLICIT REAL*8 (A-H,O,P,R-Y), COMPLEX*16 (C,Z)
DIMENSION Q1ABCF(13),Q1RBCF(13)
IF(IPT,GT,0) WRITE(6,990) DK,DR,DW,RHO,CFFSET,SIGMA,XSTN,IPT
FORMAT(0,10X,'Q1OPHI ENTERED WITH:',/
1 10X,'DK:',11X,'DR:',11X,'DW:',11X,'RHO:',10X,'OFFSET',7X,'SIGMA',
2 10X,'XSTN:',9X,'IPT',/,',10X,7(E12.5,',' ,13)
RZERO=0.0D0
ICT=IPT-1
C1OPHI=Q1CRBP(DK,DR,DW,RZERO,XSTN,IOT)
1 + Q1DRIP(DK,DR,DW,RZERO,OFFSET,SIGMA,XSTN,IOT)
2 + Q1DABP(DK,DR,DW,DR,CFFSET,SIGMA,XSTN,IOT)
3 + Q1DAIP(DK,DR,DW,DR,OFFSET,SIGMA,XSTN,Q1ABCF,N,IOT)
IF(IPT,LE,0) RETURN
WRITE (6,995) Q1OPHI
995 FCRMAT(0,10X,'C1OPHI = ',E14.7,',' ,E14.7)
RETURN
ENC

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CCMPLEX FUNCTION C1CREP*16(CK,DR,CLW,RHO,XSTN,IPT)
IMPLICIT REAL*8 (A-H,C,P,R-Y), COMPLEX*16 (Q,Z)
DIMENSION CIMP(2)
IF (IPT.GT.0) WRITE(6,990) CK,DR,CLW,RHO,XSTN,IPT
FCRMT(C,10X,'QICRBP ENTERED WITH: ',/,
1,10X,10X,CLW,DR,11X,DW,11X,RHC,10X,'XSTN',9X,'IPT',
2,10X,5(E12.5),11X,13)
IF(XSTN.LE.RHO-1.000) GOTO 20
ICT=IPT-1
GAMMA=1.400
DM2=(GAMMA+1.000)*DW
CLAMDA=(DK/DM2)
DM2=(GAMMA+1.000)*DW
PIMAG=(DK+DLAMDA)*(XSTN-RHO)
QICRBP=-DCMPLX(1.000,PIMAG-CLAMDA*XSTN)/DSQRT(DM2)
IF(IPT.LE.0) RETURN
GOTO 60
QICRBP=DCMPLX(0.000,0.000)
20 IF(IPT.LE.0) RETURN
60 WRITE(6,995) QIDRBP
995 RETURN
END
CCMPLEX FUNCTION C1DABP*16(DK,DR,CLW,RHO,OFFSET,SIGMA,XSTN,IPT)
IMPLICIT REAL*8 (A-H,C,P,R-Y), COMPLEX*16 (Q,Z)
DIMENSION CIMP(2)
IF (IPT.GT.0) WRITE(6,990) DK,DR,CLW,RHO,OFFSET,SIGMA,XSTN,IPT
FCRMT(C,10X,10X,QIDRBP ENTERED WITH: ',/,
1,10X,10X,CLW,DR,11X,DW,11X,RHC,10X,'XSTN',9X,'IPT',
2,10X,5(E12.5),11X,13)
IF(XSTN.LE.RHO-1.000) GOTO 20
ICT=IPT-1
GAMMA=1.400
DM2=(GAMMA+1.000)*DW
CLAMDA=(DK/DM2)
PIMAG=(DK+DLAMDA)*(XSTN-RHO)
QICRBP=-DCMPLX(1.000,PIMAG-CLAMDA*XSTN)*CCNST/CSQRT(DM2)
IF(IPT.LE.0) RETURN
GOTO 60
QICRBP=DCMPLX(0.000,0.000)
20 IF(IPT.LE.0) RETURN
60 WRITE(6,995) QIDRBP
995 RETURN
END
CCMPLEX FUNCTION C1PAIP*16(CK,DR,CLW,RHO,CFST,SGMA,XSTN,CICF,N,IPT)

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